

mechanical models for second law violation

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outline

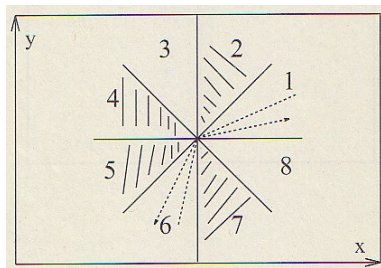
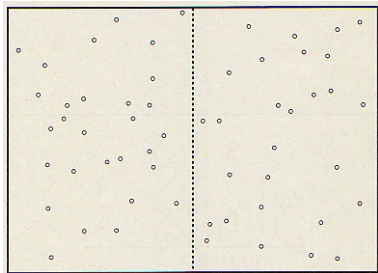
- ▶ the goal is to deduce the properties of mechanical models violating the second law

- ▶ finally we will see that:
violation of the second law requires noninvariance of phase-space volume

model 1:

P.A. Skordos: Compressible dynamics, time reversibility, Maxwell's demon, and the second law

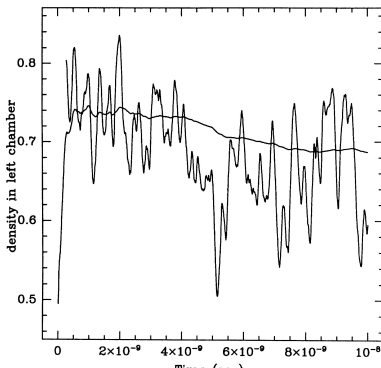
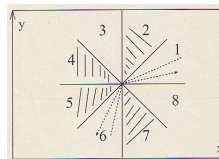
- ▶ operation is time-reversible
- ▶ speed and kinetic energy are conserved
- ▶ membrane more permeable from right than from left \Rightarrow density difference
- ▶ equilibrium reached irreversibly when fluxes from both sides become equal



Estimate of density difference

- ▶ Assuming that the velocities of the disks are distributed isotropically \Rightarrow (impact rate) $\propto |\cos \vartheta|$, ϑ =the impact angle

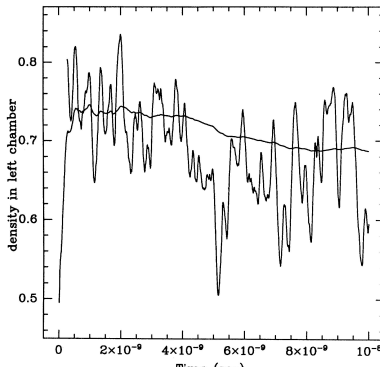
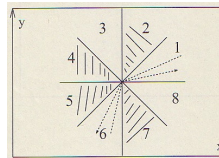
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Compression of phase-space volume I

1. consider only one disk, characterized by (x, y, ϑ) . Look at an infinitesimal phase volume $\omega(x_1, y_1, \vartheta_1)$ centred at a the point (x_1, y_1, ϑ_1) .

Assume: membrane at $x = 0$, $x_1 > 0$ near $x = 0$, $\vartheta_1 \in (\frac{3\pi}{4}, \pi)$.

Also assume that after a time interval 1 that every point in $\omega(x_1, y_1, \vartheta_1)$ has moved to the left side of the membrane.

2. let $(\tilde{x}, \tilde{y}, \tilde{\vartheta})$ denote the image of the evolution map

$$\tilde{x} = -\sin \vartheta - x \frac{\sin \vartheta}{\cos \vartheta}$$

$$\tilde{y} = y - x \left(1 + \frac{\sin \vartheta}{\cos \vartheta} \right) - \cos \vartheta$$

$$\tilde{\vartheta} = \frac{3\pi}{2} - \vartheta$$

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other membrane maps

Is there another membrane map that would result in a density difference while preserving phase space volume?

1. We seek: $f : \vartheta \rightarrow f(\vartheta)$ s. t. the Jacobian determinant equals 1.
 $\Rightarrow f(\vartheta) = \arcsin(\pm \sin \vartheta + C) \quad C = \text{constant}$
So the only thing left to do is to find a function $f(\vartheta)$, which fulfills $P_{(L \rightarrow R)} \neq P_{(R \rightarrow L)}$

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Therefore a membrane map (piecewise differentiable) that obeys incompressible dynamics cannot create a density difference.

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Maxwell's demon

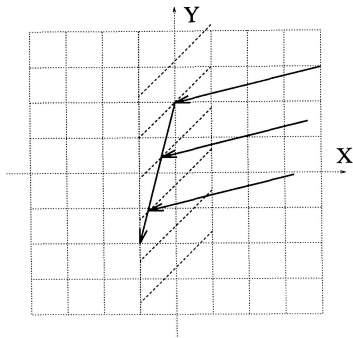
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Maxwell's
demon:(time-)irreversible
- ▶ membrane \longleftrightarrow tennis demon.
but here trajectories evolve
bijectively, so no info erasure.
How to nullify Maxwell's demon
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locations, all separated by Δx
 \Rightarrow many-to-one map

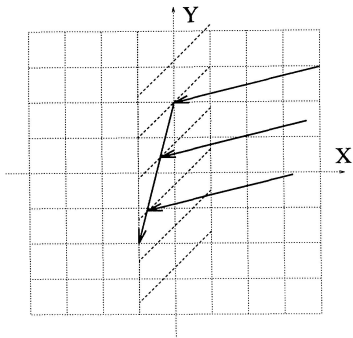
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model 2:

Kechen and Kezhao Zhang: Mechanical models of Maxwell's demon with noninvariant phase volume

mechanical models of Maxwell's demon I

- ▶ our starting point is a system of N interacting point particles:

$$m_i \ddot{\vec{r}}_i = \vec{F}(\vec{r}_1, \dots, \vec{r}_N, \dot{\vec{r}}_1, \dots, \dot{\vec{r}}_N)$$

- ▶ we will use velocity-dependent force fields to implement a barrier which allows particles to pass through preferably only in one direction. We also demand: $\vec{F} \perp \vec{v} \Rightarrow$ the total energy is conserved.

$$\vec{F}(\vec{v}) = (\vec{v} \times \hat{z}) A(v) B(\vartheta)$$

- ▶ $\vec{F}(\vec{v}) = (\vec{v} \times \hat{z}) A(v) B(\vartheta)$ is time-reversible $\Leftrightarrow \vec{F}(\vec{v}) = \vec{F}(-\vec{v})$
- ▶ phase volume conserved $\Leftrightarrow B(\vartheta) = \text{const}$

mechanical models of Maxwell's demon II

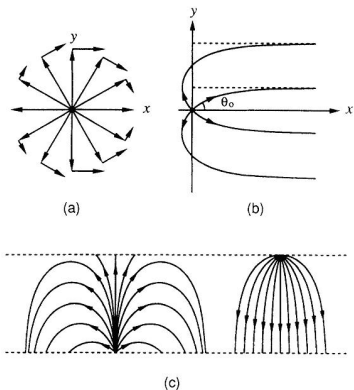


Figure: c) uniform field between dashed lines, downward: attracting direction, leakage caused by nearly perpendicular impinging

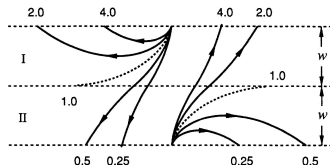


Figure: in I: attracting direction is upwards; in II: attracting direction is downwards; barrier width w chosen such that particles with speed $v < 1$ pass downward, $v > 1$ pass upward

is noninvariance of phase volume necessary for Maxwell's demon I

- ▶ let us introduce SMF(spontaneous momentum flow): sustaining and robust momentum flow inside an isolated mechanical system.
- ▶ there is a SMF in the system if the long-term average $\bar{J}_\nu = \lim_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau J_\nu(T_t \zeta) dt \neq 0 \quad \forall \zeta \in \Sigma (= \text{energy surface})$ for some spatial region V

(where $J_\nu(\zeta) = \text{total momentum inside } V \text{ at a given time} = \sum_{i=1}^N p_i \chi_V(r_i)$)

is noninvariance of phase volume necessary for Maxwell's demon II

Formulation of the second law by SMF

existence of perpetual motion machine of the second kind \Leftrightarrow
existence of SMF

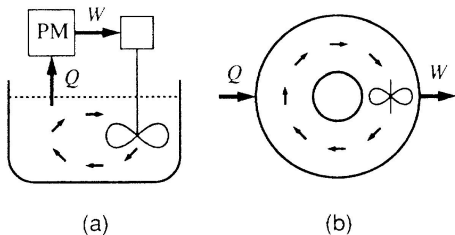


Figure: (a) \Rightarrow (b) \Leftarrow

is noninvariance of phase volume necessary for Maxwell's demon III

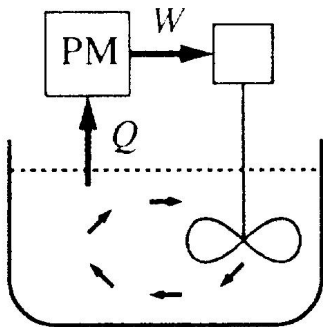
Theorem: nonexistence of SMF

The point particle system $m_i \ddot{\vec{r}}_i = \vec{F}(\vec{r}_1, \dots, \vec{r}_N, \dot{\vec{r}}_1, \dots, \dot{\vec{r}}_N)$ **cannot** exhibit SMF if

- i) its energy function E is symmetric under momentum reversal, namely, $E(\zeta) = E(\tilde{\zeta})$ where $\zeta = (q_1, \dots, q_s, p_1, \dots, p_s)$ and $\tilde{\zeta} = (q_1, \dots, q_s, -p_1, \dots, -p_s)$
- ii) phase volume $dq_1 \cdots dq_s dp_1 \cdots dp_s$ is invariant during time evolution, and total phase volume is finite for finite energy

is noninvariance of phase volume necessary for Maxwell's demon IV

Theorem \Rightarrow such a system cannot serve as Maxwell's demon,
because: suppose we can \rightarrow contradiction to theorem



summary and conclusion

model 1

- ▶ time-reversible microscopic dynamics
- ▶ violation of second law \Rightarrow compression of phase space

model 2

- ▶ time-reversible microscopic dynamics and purely mechanical
- ▶ In systems with symmetric energy wrt momentum reversal, 2nd law violation requires noninvariance of phase volume.

ultimate result

The invariance of phase volume appears as a factor responsible for the validity of the second law of thermodynamics