

Exercise 10.1 Uniaxial Compressibility

We consider a system of electrons upon which an uniaxial pressure in z-direction acts. Assume that this pressure causes a deformation of the Fermi surface $k \equiv k_F^0$ of the form

$$k_F(\phi, \theta) = k_F^0 + \gamma \frac{1}{k_F^0} \left[3k_z^2 - (k_F^0)^2 \right] = k_F^0 + \gamma k_F^0 [3 \cos^2 \theta - 1], \quad (1)$$

where $\gamma = (P_z - P_0)/P_0$ is the anisotropy of the applied pressure.

- Show that for small $\gamma \ll 1$, the deformed Fermi surface $k_F(\phi, \theta)$ encloses the same volume as the non-deformed one, k_F^0 , where terms of order $\mathcal{O}(\gamma^2)$ can be neglected.
- The deformation of the Fermi surface effects a change in the distribution function of the electrons. Using Landau's Fermi Liquid theory, calculate the uniaxial compressibility

$$\kappa_u = \frac{1}{V} \frac{\partial^2 E}{\partial P_z^2}, \quad (2)$$

which is caused by the deformation given in eq. (1) (E denotes the Landau energy functional).

Exercise 10.2 Polarization of a neutral Fermi liquid

Consider a system of neutral spin-1/2 particles each carrying a magnetic moment $\boldsymbol{\mu} = \frac{\mu}{2} \boldsymbol{\sigma}$. An electric field \mathbf{E} couples to the atoms by the relativistic spin-orbit interaction

$$H_{SO} = \frac{\mu}{2} \left(\frac{\mathbf{v}}{c} \times \mathbf{E} \right) \cdot \boldsymbol{\sigma} \quad (3)$$

where $\boldsymbol{\sigma} = (\sigma^x, \sigma^y, \sigma^z)$ is the vector of Pauli spin matrices. In the following we want to calculate the linear response function χ for the uniform polarization

$$\mathbf{P} = \chi \mathbf{E}. \quad (4)$$

In the presence of spin-orbit interaction we have to consider a more general situation of a distribution of quasiparticles with variable spin quantization axis. In such a case we must treat the quasiparticle distribution function and the energy as a 2×2 matrix, $(\hat{n}_{\mathbf{p}})_{\alpha\beta}$ and $(\hat{\epsilon}_{\mathbf{p}})_{\alpha\beta}$, respectively. Furthermore, we require f to be a scalar under spin rotations. In this case f must be of the form

$$\hat{f}_{\alpha\beta, \alpha'\beta'}(\mathbf{p}, \mathbf{p}') = f^s(\mathbf{p}, \mathbf{p}') \delta_{\alpha\beta} \delta_{\alpha'\beta'} + f^a(\mathbf{p}, \mathbf{p}') \boldsymbol{\sigma}_{\alpha\beta} \cdot \boldsymbol{\sigma}_{\alpha'\beta'} \quad (5)$$

- Expand $\hat{n}_{\mathbf{p}}$, $\hat{\epsilon}_{\mathbf{p}}$, and $\hat{f}_{\boldsymbol{\sigma}\boldsymbol{\sigma}'}(\mathbf{p}, \mathbf{p}')$ in terms of the unit matrix $\sigma^0 = \mathbf{1}$ and the Pauli spin matrices $\sigma^1 = \sigma^x$, $\sigma^2 = \sigma^y$, $\sigma^3 = \sigma^z$ and find Landau's energy functional E .
- Assume that the electric field is directed along the z direction. Show that the polarization of such a system is given by

$$P_z = \frac{\partial E}{\partial E_z} = \frac{\mu}{m^* c} \sum_{\mathbf{p}} (p_y \delta n_{\mathbf{p}}^1 - p_x \delta n_{\mathbf{p}}^2). \quad (6)$$

Here, $\delta n_{\mathbf{p}}^i = \frac{1}{2} \text{tr} [\delta \hat{n}_{\mathbf{p}} \sigma^i]$ and $\delta \hat{n}_{\mathbf{p}}$ is the deviation from the equilibrium ($E_z = 0$) distribution function.

- c) The application of the electric field changes the quasiparticle energy in linear response according to

$$\delta\tilde{\epsilon}_{\mathbf{p}}^i = \delta\epsilon_{\mathbf{p}}^i + \frac{2}{V} \sum_{\mathbf{p}'} f^{ii}(\mathbf{p}, \mathbf{p}') \delta n_{\mathbf{p}'}^i \quad \text{with} \quad \delta n_{\mathbf{p}}^i = \frac{\partial n_0}{\partial \epsilon} \delta\tilde{\epsilon}_{\mathbf{p}}^i = -\delta(\epsilon_{\mathbf{p}}^0 - \epsilon_F) \delta\tilde{\epsilon}_{\mathbf{p}}^i. \quad (7)$$

Use the ansatz $\delta\tilde{\epsilon}_{\mathbf{p}}^i = \alpha \delta\epsilon_{\mathbf{p}}^i$ and show that $\alpha = 1/(1 + F_1^a/3)$ to find $\delta n_{\mathbf{p}}^i$ and $\delta\tilde{\epsilon}_{\mathbf{p}}^i$.

- d) Compute χ according to Eq. (4).

Office hour:

Monday, May 9th, 2011 - 9:00 to 11:00 am

HIT K 12.1

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