

**Exercise 1. Optical Theorem**

Writing the relation between the  $S$ -matrix and the  $T$ -matrix as

$$S_{\beta\alpha} = \delta(\beta - \alpha) - i 2\pi \delta(E_\alpha - E_\beta) T_{\beta\alpha}, \quad (1)$$

the unitarity condition on the  $S$ -matrix reads

$$(S^\dagger S)_{\beta\alpha} = \sum_{\gamma} (S_{\gamma\beta})^* (S_{\gamma\alpha}) = \delta(\beta - \alpha),$$

with a sum over all intermediate states  $\gamma$  including an integration over their momenta.

(a) Show that the unitarity condition for  $\beta = \alpha$  implies

$$-2 \operatorname{Im}(T_{\alpha\alpha}) = (2\pi) \sum_{\gamma} \delta(E_\gamma - E_\alpha) |T_{\gamma\alpha}|^2 \quad (2)$$

(b) Starting from the equation above, prove the Optical Theorem:

$$\sigma_{\text{tot}} = \frac{4\pi}{k_\alpha} \operatorname{Im}(f_{\alpha\alpha}) \quad (3)$$

where  $f_{\alpha\alpha} = -\frac{m_\alpha}{2\pi\hbar^2} T_{\alpha\alpha}$  is the elastic scattering amplitude in the forward direction.

*Hint.* Recall that the differential cross-section for the process  $\alpha \rightarrow \gamma$  can be written as:

$$\frac{d\sigma_{\alpha \rightarrow \gamma}}{d\Omega} = \frac{1}{\text{flux}} \frac{2\pi}{\hbar} |T_{\gamma\alpha}|^2 \frac{d\rho}{d\Omega},$$

where  $\rho$  is the density of states

$$\rho = \int \frac{d^3\vec{k}_\gamma}{(2\pi)^3} \delta(E_\gamma - E_\alpha)$$

**Exercise 2. The photoelectric effect**

We want to treat the photoelectric effect using the quantization of the electromagnetic field. A photon is absorbed by the atom and an electron is excited from the bound state:

$$\psi_i(\vec{r}) = \sqrt{\frac{Z^3}{\pi a_0^3}} e^{-\frac{Zr}{a_0}}, \quad (4)$$

where  $a_0 = \frac{\hbar^2}{m_e e^2}$ , to a free state

$$\psi_f(\vec{r}) = e^{i\vec{k}_e \cdot \vec{r}}. \quad (5)$$

(a) Compute the transition matrix element  $V_{fi} = \langle \psi_f; (n-1)(\vec{k}_\lambda, \lambda) | V | \psi_i; n(\vec{k}_\lambda, \lambda) \rangle$ .

- (b) Using the Fermi's golden rule for a continuum of states, show that the differential cross section is given by:

$$\frac{d\sigma}{d\Omega} = \frac{e^2}{\hbar c} \frac{32 \hbar k_e}{m \omega_\lambda} (\vec{\varepsilon}_\lambda \cdot \vec{k}_e)^2 \frac{\left(\frac{Z}{a_0}\right)^5}{\left[\left(\frac{Z}{a_0}\right)^2 + |\vec{k}_\lambda - \vec{k}_e|^2\right]^2}, \quad (6)$$

where  $\vec{\varepsilon}_\lambda, \vec{k}_\lambda, \omega_\lambda$  corresponds to the absorbed photon described by  $(\vec{k}_\lambda, \lambda)$ .

- (c) If the photon energy is large compared to the binding energy of the electron but small compared with the rest mass energy ( $W = \frac{mZ^2e^4}{2\hbar^2} = \frac{\hbar^2 Z^2}{2m a_0^2} \ll \hbar\omega_\lambda \ll mc^2$ ), show that for an unpolarized light (equal polarization along  $x$  and  $y$  direction) the scattering cross section can be written as

$$\frac{d\sigma}{d\Omega} = \frac{e^2}{\hbar c} 32 \left(\frac{k_e a_0}{Z}\right)^3 \left(\frac{W}{\hbar\omega_\lambda}\right)^5 \frac{\left(\frac{a_0}{Z}\right)^2 \sin^2 \theta}{\left(1 - \frac{v}{c} \cos \theta\right)^4}, \quad (7)$$

where  $v = \frac{\hbar k_e}{m}$  is the electron velocity, the  $z$ -axis coincides with the direction of the incident photon and the photon velocity is described by the polar angle  $\theta$  ( $\vec{k}_\lambda \cdot \vec{k}_e = k_\lambda k_e \cos \theta$ ).