

R. Feynman
1918-1988
Nobel laureate 1965



Tests of QED

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How to test QED theory

- Precision tests of the **Quantum Electrodynamics** theory usually consist in the measurement of the electromagnetic fine structure constant α in different systems. Experimental results are compared with theoretical predictions
- The validation process requires **very high precision** in both measurements and theoretical calculations
- QED is then confirmed to the extent that these measurements of α from different physical sources agree with each other
- The most stringent test of QED is given by the measurement of the **electron magnetic moment**. However, several other experimental tests have been performed in different energy ranges and systems:
 - ◆ Low energy range, accessible with small experiments
 - ◆ High energy range, accessible with particle colliders (e.g. e^+e^- colliders)
 - ◆ Condensed matter systems
- As we will see, the achieved precision makes QED one of the **most accurate physical theories** constructed so far

Part one:
Tests of QED at high energy

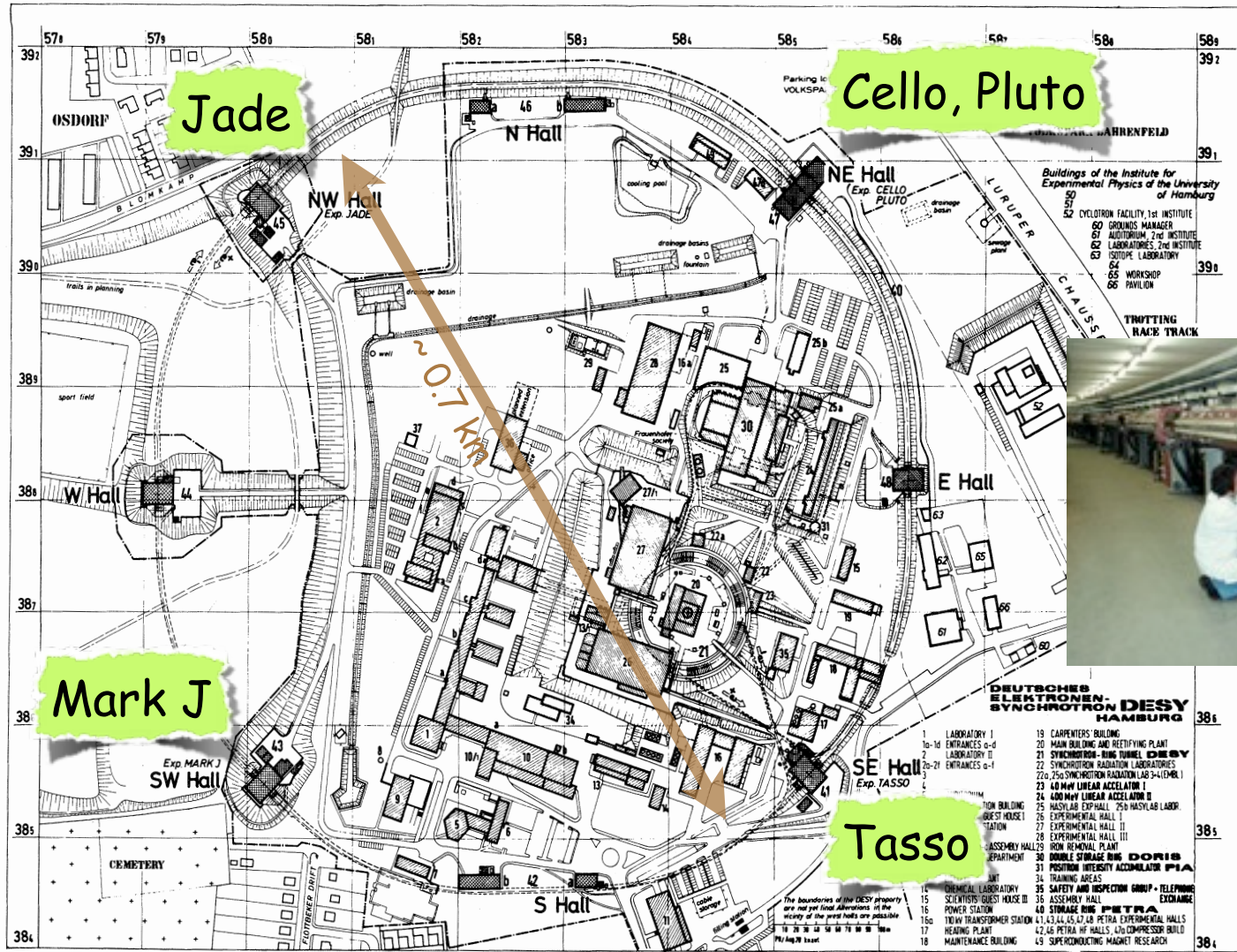
Tests of QED at high energy

- In addition to the low-energy experiments, QED has been tested also in high energy electron-positron collisions
- We discuss here the following reactions:
 - ◆ $e^+e^- \rightarrow e^+e^-$ (*Bhabha scattering*)
 - ◆ $e^+e^- \rightarrow \mu^+\mu^-, \tau^+\tau^-$ (*Lepton pair production*)
 - ◆ $e^+e^- \rightarrow \gamma \gamma$ (*Two photon annihilation*)
 - ◆ $e^+e^- \rightarrow q \bar{q} \rightarrow \text{hadrons}$ (*Total hadronic cross section*)
- The energy range (\sqrt{s}) between 12 GeV and 47 GeV was investigated with the PETRA accelerator at DESY (Hamburg)
- High energy ranges (90-200 GeV) were covered by the LEP collider at CERN (Geneva). However, electro-weak contributions to the cross-sections become considerable at these energies
- Intermediate energies were covered by TRISTAN and SLC

e+e- colliders

Accelerator	Experiments	CMS-Energy	Integrated luminosity
SPEAR	SPEAR	2 - 8 GeV	—
PEP	ASP, DELCO, HRS, MARK II, MAC	up to 29 GeV	~ 300 pb ⁻¹
PETRA	JADE, MARK J PLUTO, TASSO CELLO	12 - 47 GeV	~ 20 pb ⁻¹
TRISTAN	TRISTAN	50 - 60 GeV	~ 20 pb ⁻¹
SLC	MARK II, SLD	90 GeV	~ 25 pb ⁻¹
LEP	ALEPH, DELPHI, OPAL, L3	90 - 200 GeV	~ 200 pb ⁻¹ ~ 700 pb ⁻¹

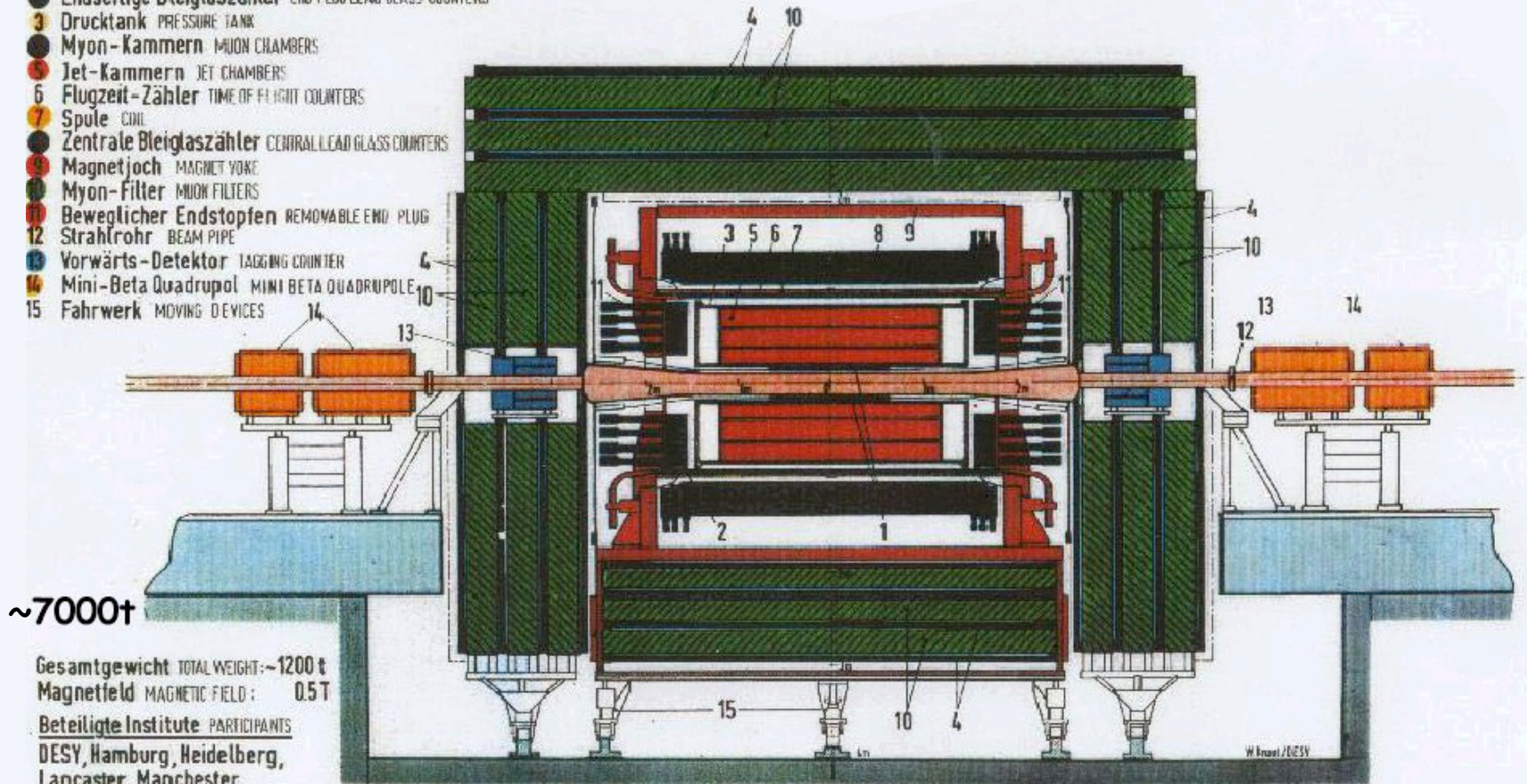
PETRA e+e- storage ring



Example: JADE detector

MAGNETDETEKTOR **JADE** MAGNET DETECTOR

- 1 Strahlrohrzähler BEAM PIPE COUNTERS
- Endseitige Bleiglaszähler END PLUG LEAD GLASS COUNTERS
- Drucktank PRESSURE TANK
- Myon-Kammern MUON CHAMBERS
- Jet-Kammern JET CHAMBERS
- 6 Flugzeit-Zähler TIME OF FLIGHT COUNTERS
- Spule COIL
- Zentrale Bleiglaszähler CENTRAL LEAD GLASS COUNTERS
- Magnetjoch MAGNET YOKE
- Myon-Filter MUON FILTERS
- Beweglicher Endstopfen REMOVABLE END PLUG
- 12 Strahlrohr BEAM PIPE
- Vorwärts-Detektor TAGGING COUNTER
- Mini-Beta Quadrupol MINI BETA QUADRUPOLE
- 15 Fahrwerk MOVING DEVICES



~7000t

Gesamtgewicht TOTAL WEIGHT: ~1200 t

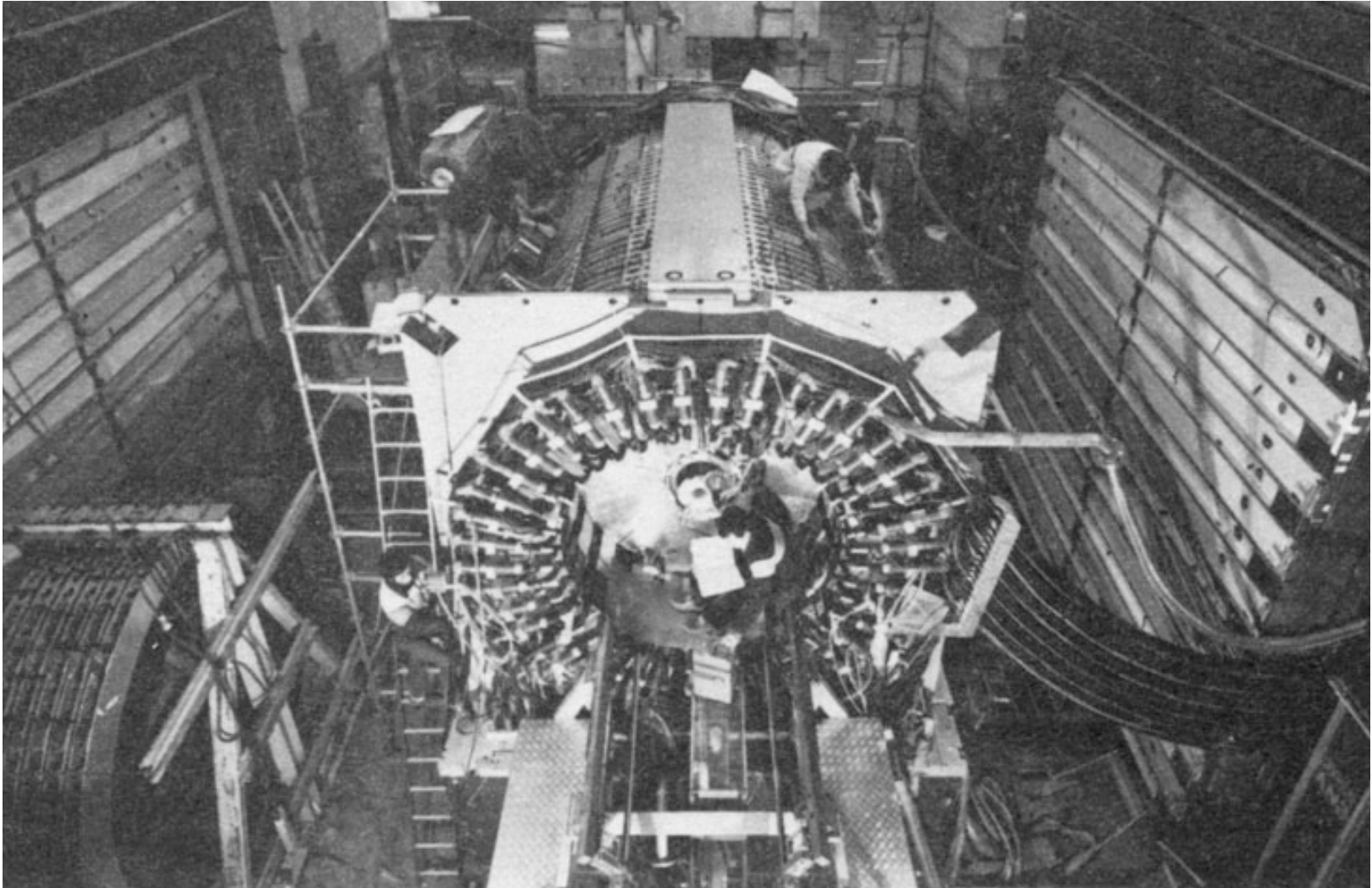
Magnetfeld MAGNETIC FIELD: 0.5T

Beteiligte Institute PARTICIPANTS

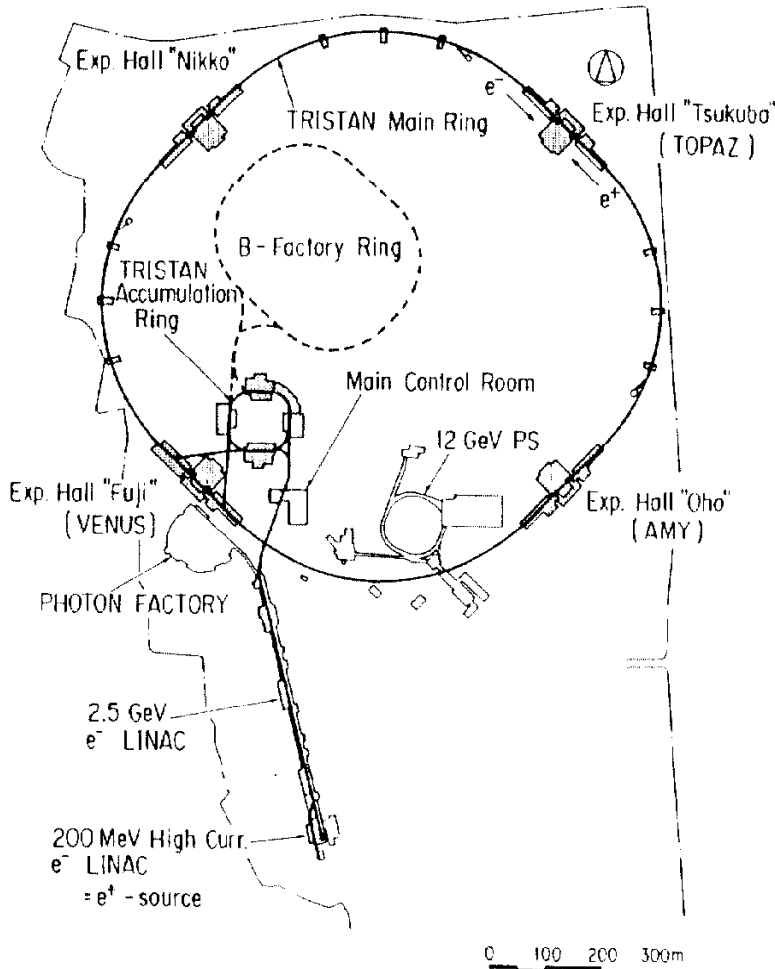
DESY, Hamburg, Heidelberg,
Lancaster, Manchester,
Rutherford Lab., Tokio

W. Krauss / DESY

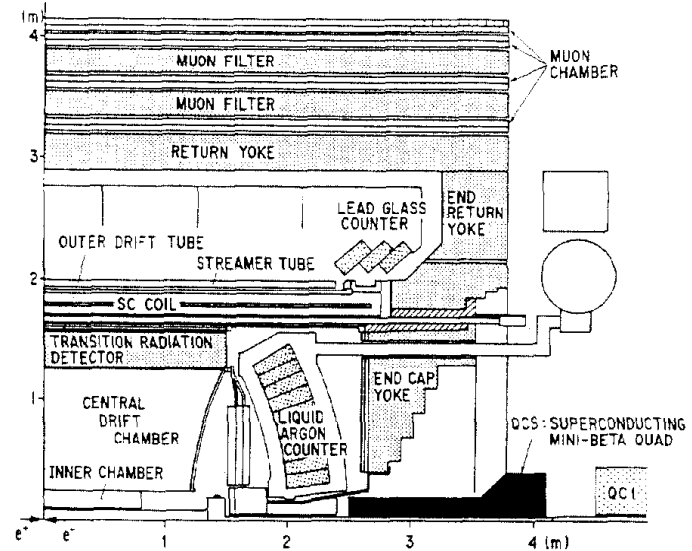
Example: JADE detector



TRISTAN accelerator & detector

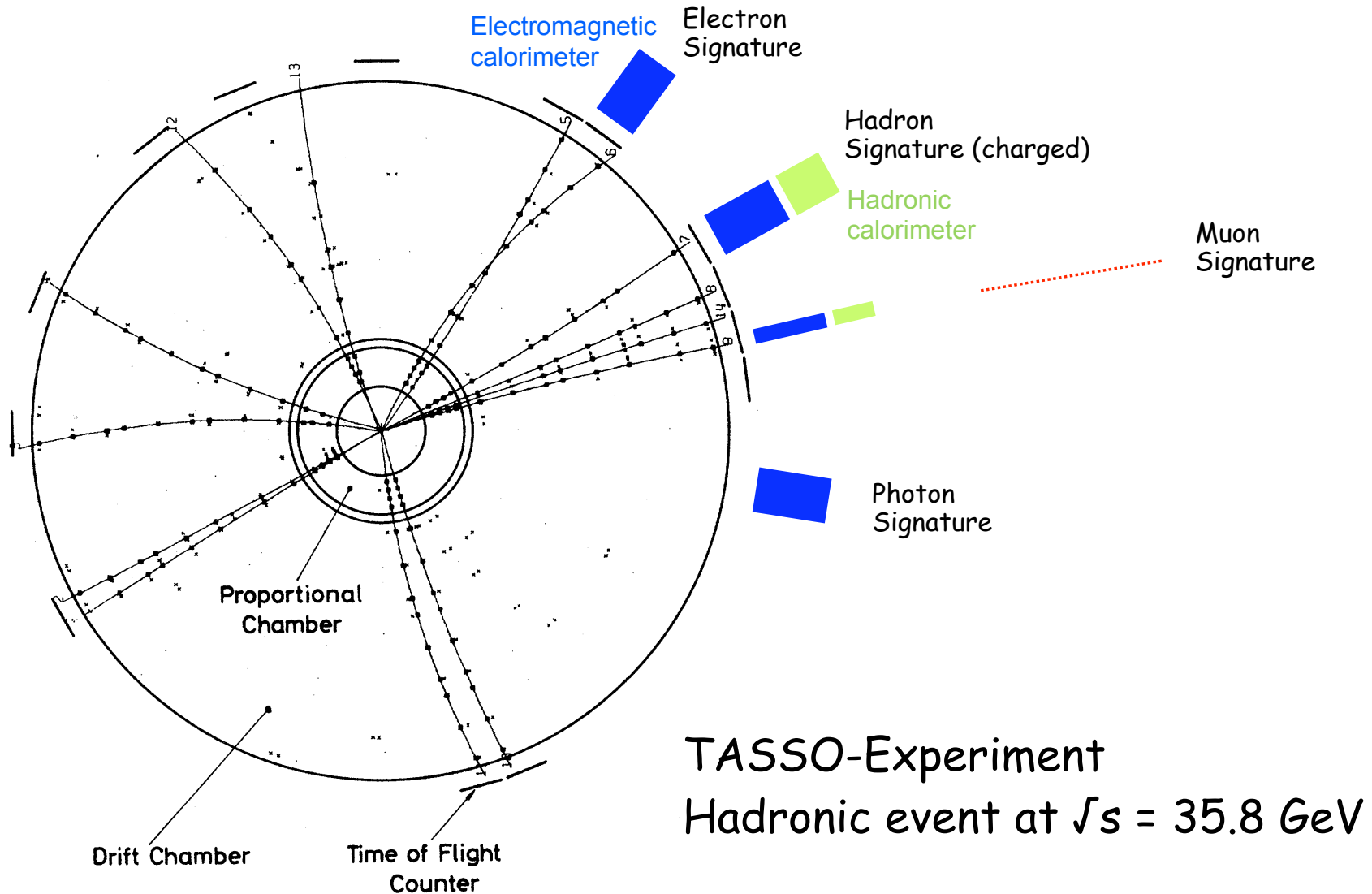


Schema of the detector

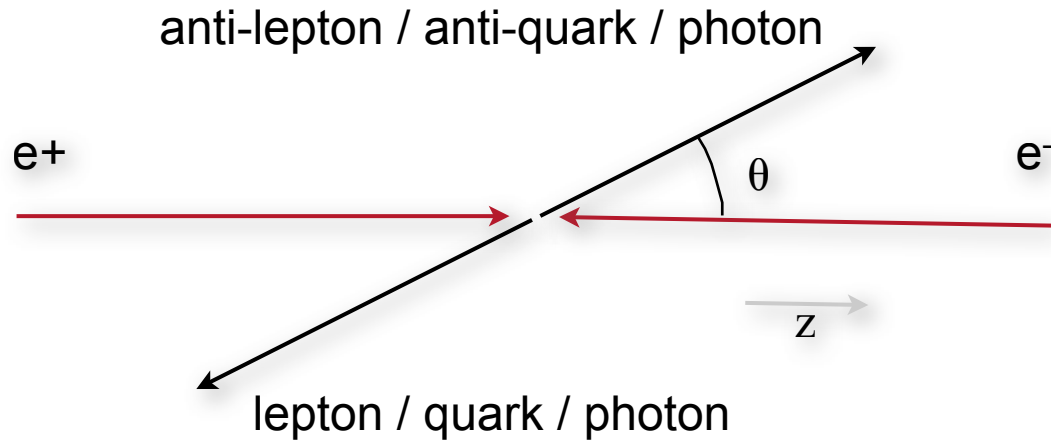


Max. beam energy: 32 GeV
 Injection energy: 8 GeV
 Beam lifetime: 5-6 hr.
 Peak luminosity: $1.4 \times 10^{31} \text{ cm}^{-2}\text{s}^{-1}$

Event reconstruction



Reminder: e^+e^- kinematics



Kinematic variables

$$q^2 = -s \frac{(1 - \cos \theta)}{2} \quad u$$

$$q'^2 = -2 \frac{(1 + \cos \theta)}{2} \quad t$$

$$\beta = \frac{p}{E}$$

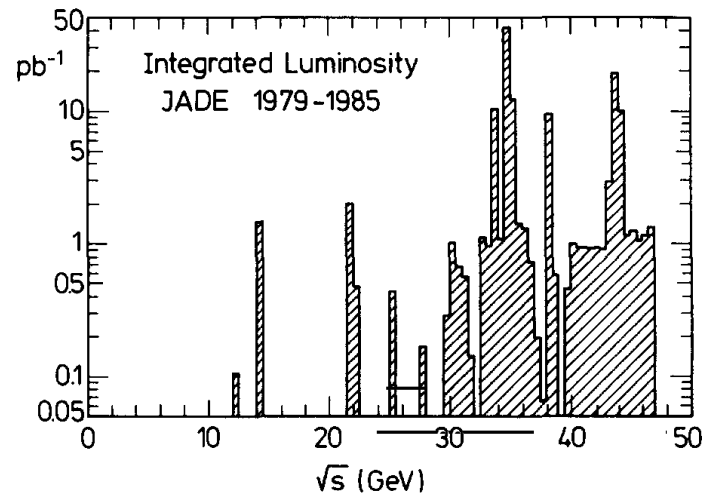
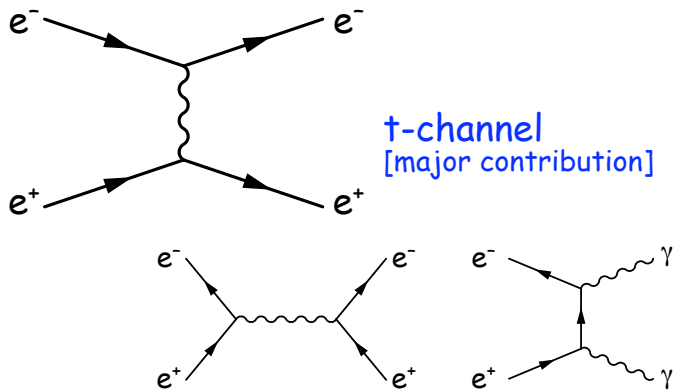
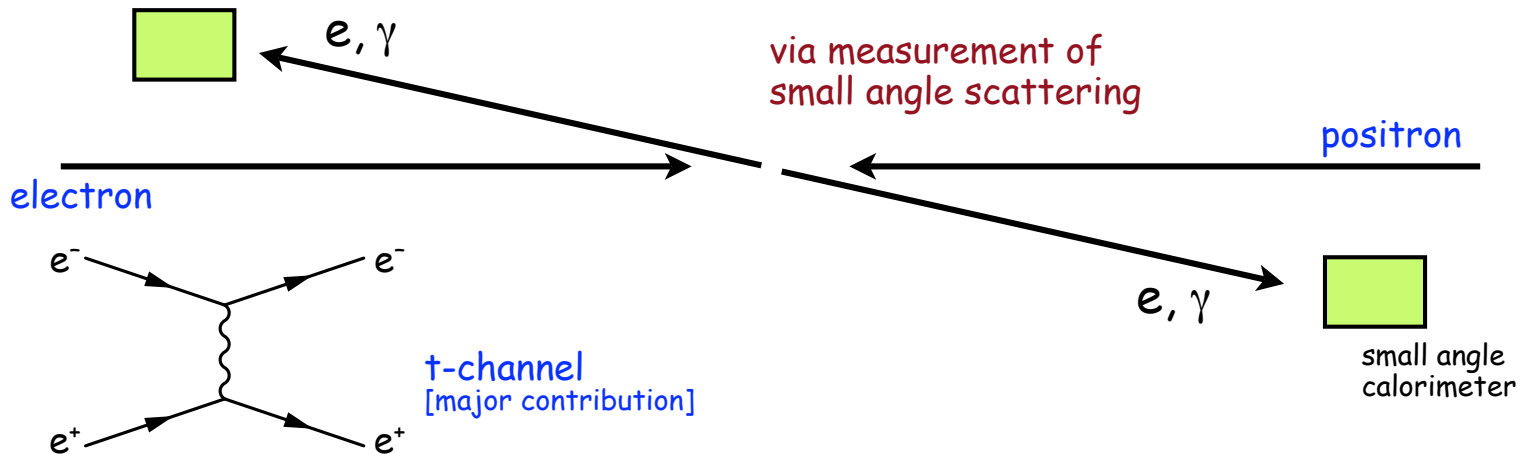
Expected cross section

$$\frac{d\sigma_{QED}}{d\Omega} = \frac{d\sigma_0}{d\Omega} (1 + \delta_{rad})$$

Radiative
(higher order)
corrections

How do we measure a cross section?

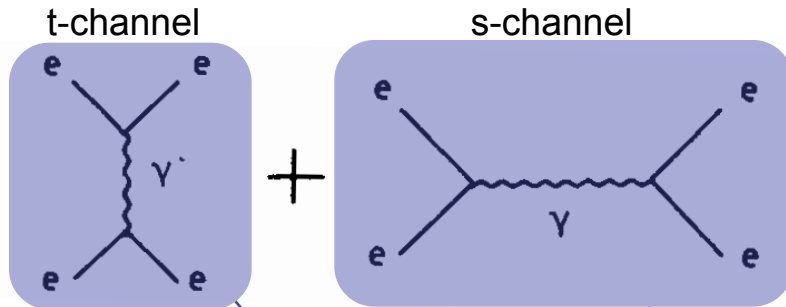
- To measure a cross section we divide the measured number of events by the integrated luminosity: $\sigma = N/L$



Calculate $\sigma_{ee,\gamma\gamma}^{\text{theo}} = \frac{N(1-b)}{(\epsilon A) \cdot L}$ Measure

Determine

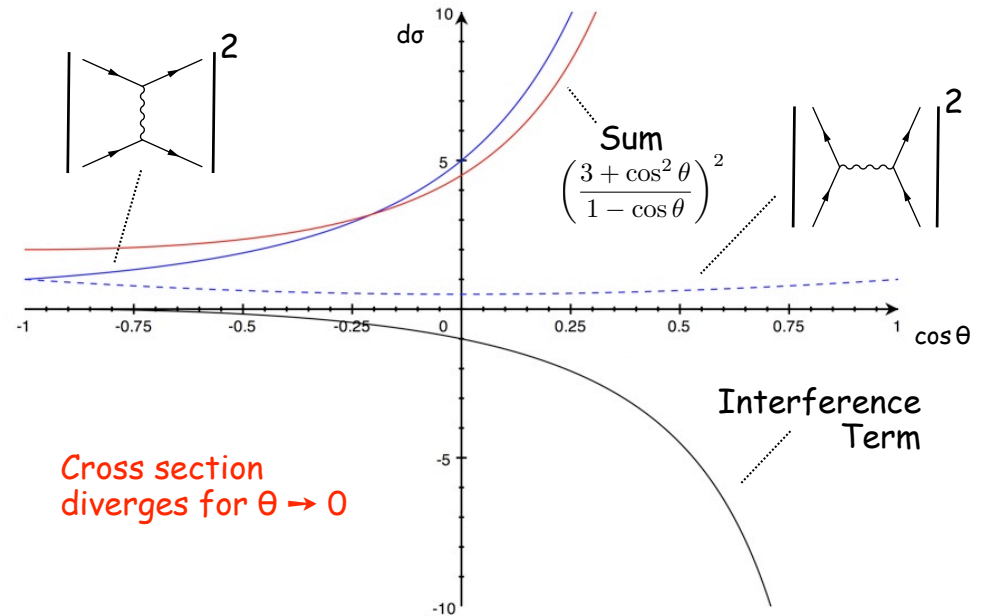
Bhabha scattering



$$\frac{d\sigma_0}{d\Omega} = \frac{\alpha^2}{2s} \left(\frac{q'^4 + s^2}{q^4} + \frac{2q'^4}{q^2s} + \frac{q'^4 + q^4}{s^2} \right)$$

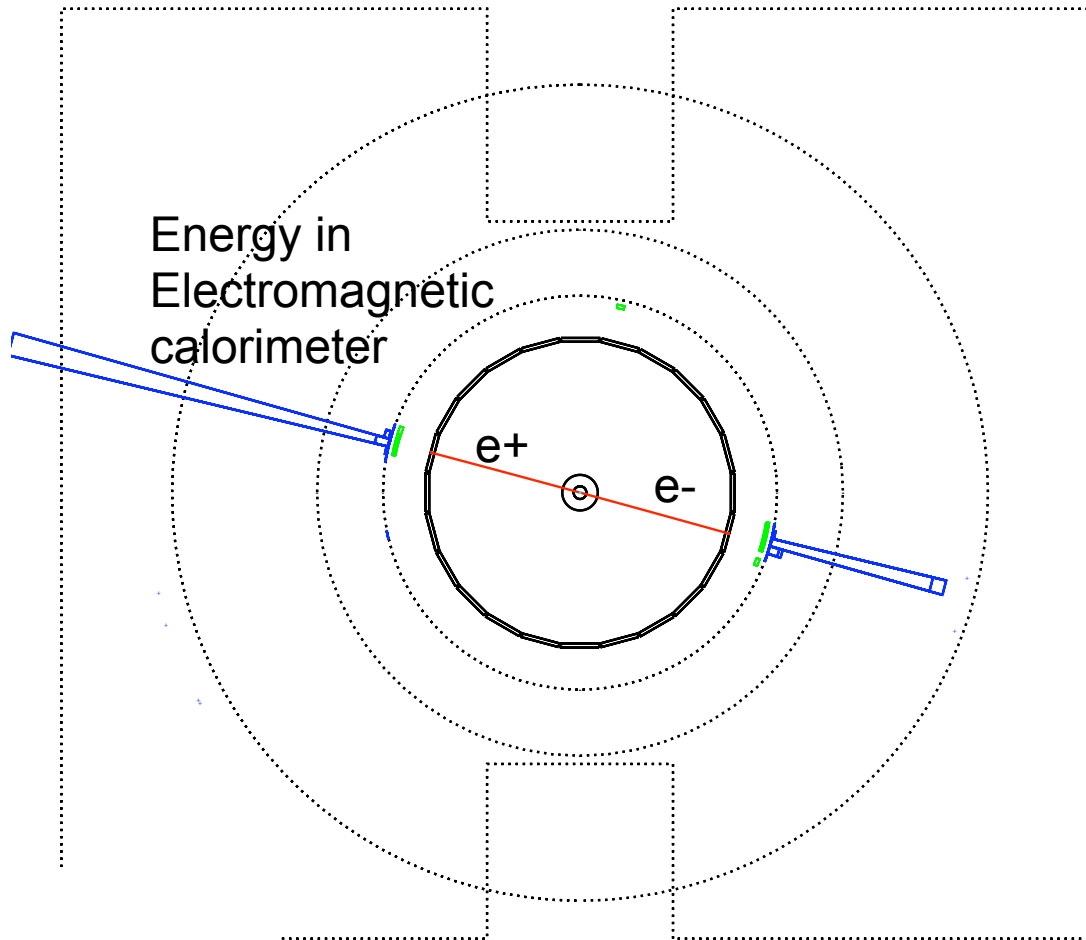
interference term

$$= \frac{\alpha^2}{4s} \left(\frac{3 + \cos^2 \theta}{1 - \cos \theta} \right)^2$$

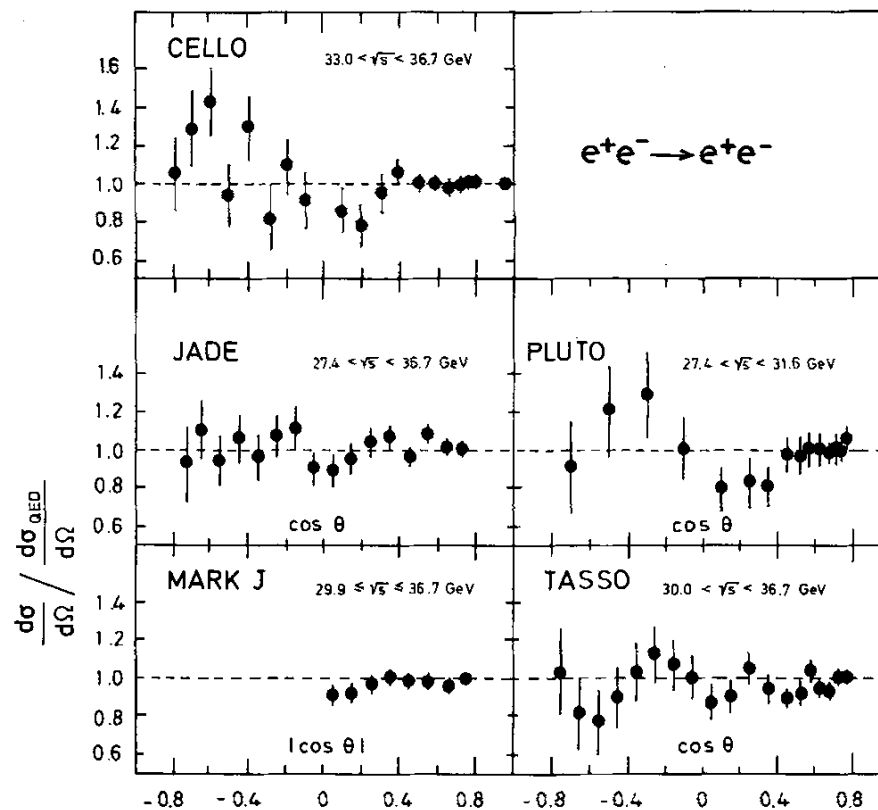
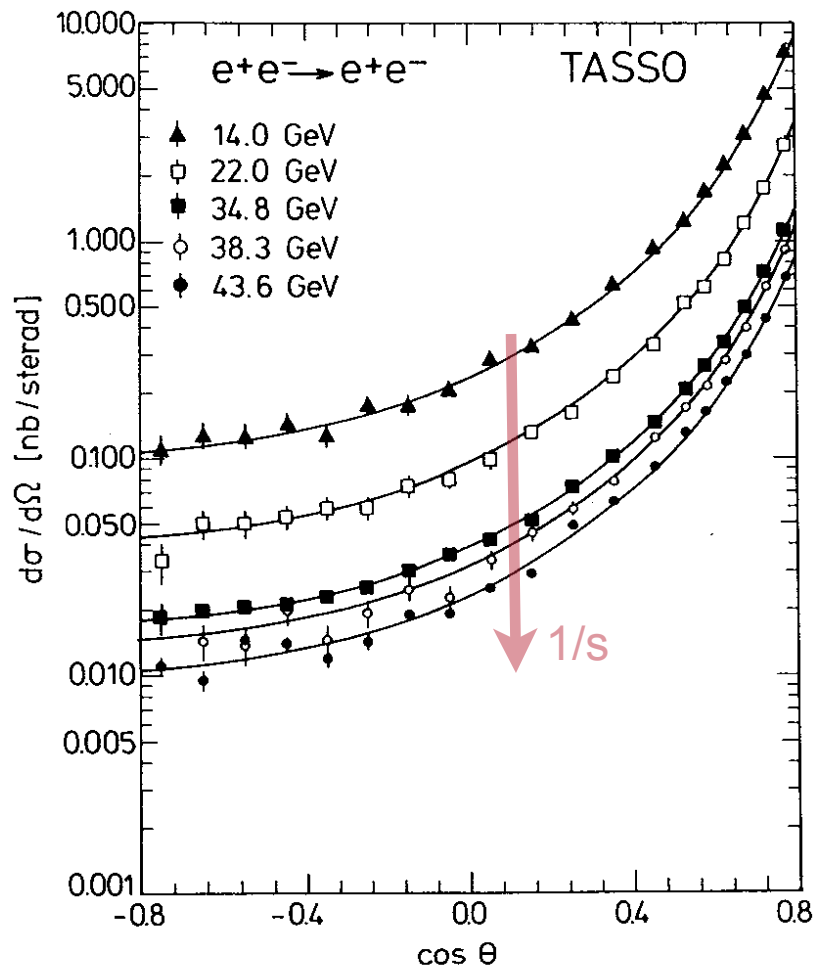


- Leading order cross section divergent for $\cos(\theta)=1$, i.e. for $\theta=0$.

Example: event display



Bhabha scattering



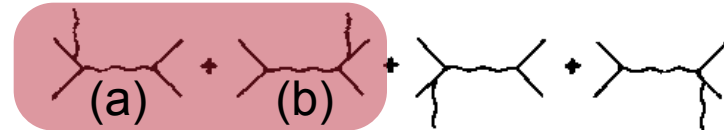
Ratio between measured cross section and QED calculation

Higher order corrections

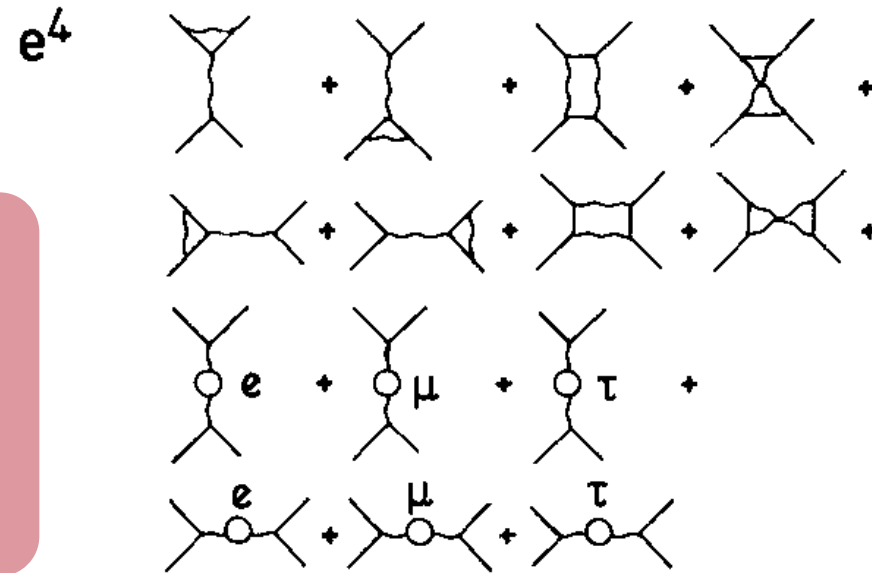
$$\propto \alpha^2 = (1/137)^2 = 5 \times 10^{-5}$$



$$\propto \alpha^3 = (1/137)^3 = 4 \times 10^{-7}$$

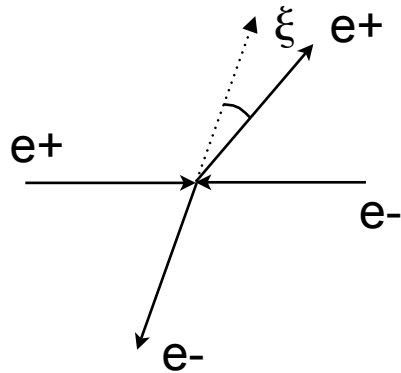


$$\propto \alpha^4 = (1/137)^4 = 2.8 \cdot 10^{-9}$$

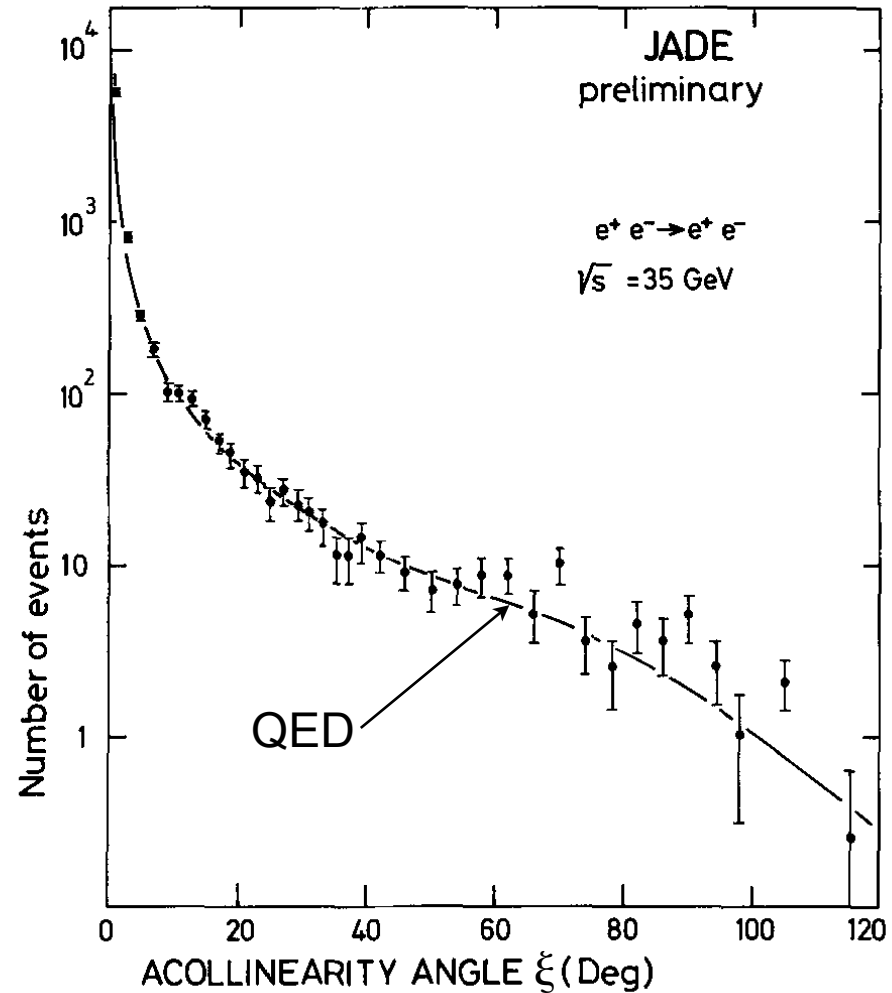


Initial (a) and final (b) state radiations produce an *acollinearity* of the final state particles:
 outgoing particles are not exactly back-to-back

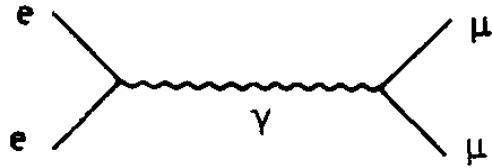
Acollinearity measurement



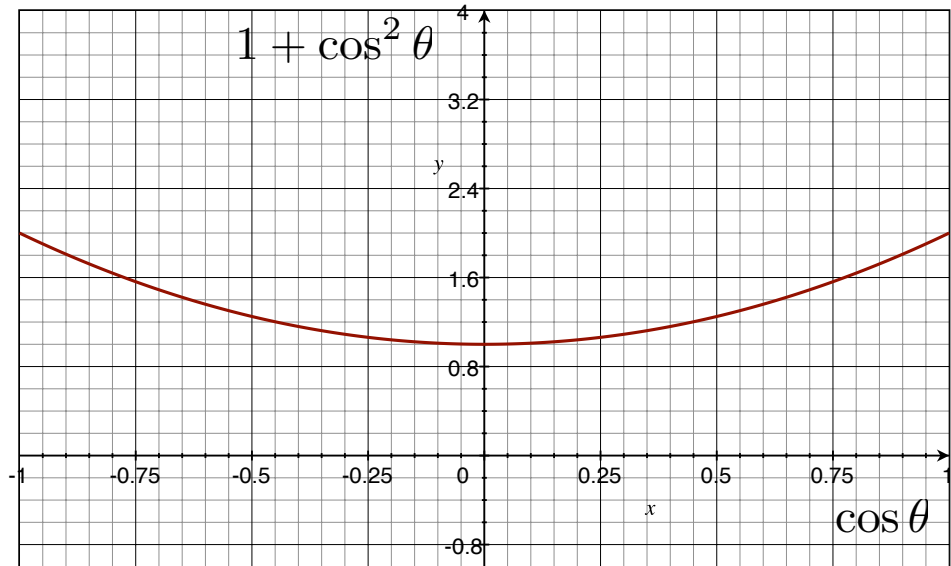
Radiative corrections
calculated at α^3 order
agree very well with the data



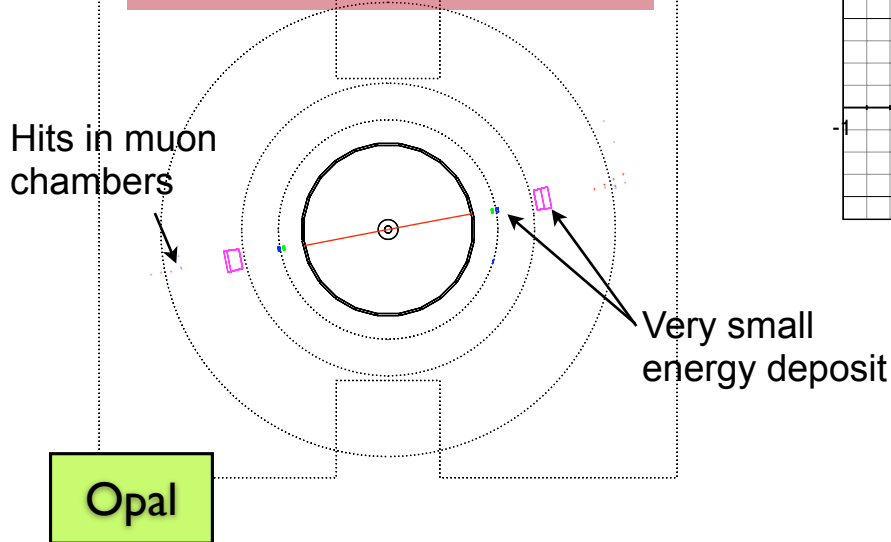
Muon pair production



$$\frac{d\sigma_0}{d\Omega} = \frac{\alpha^2}{2s} \left(\frac{q'^4 + q^4}{s^2} \right) = \frac{\alpha^2}{2s} (1 + \cos^2 \theta)$$

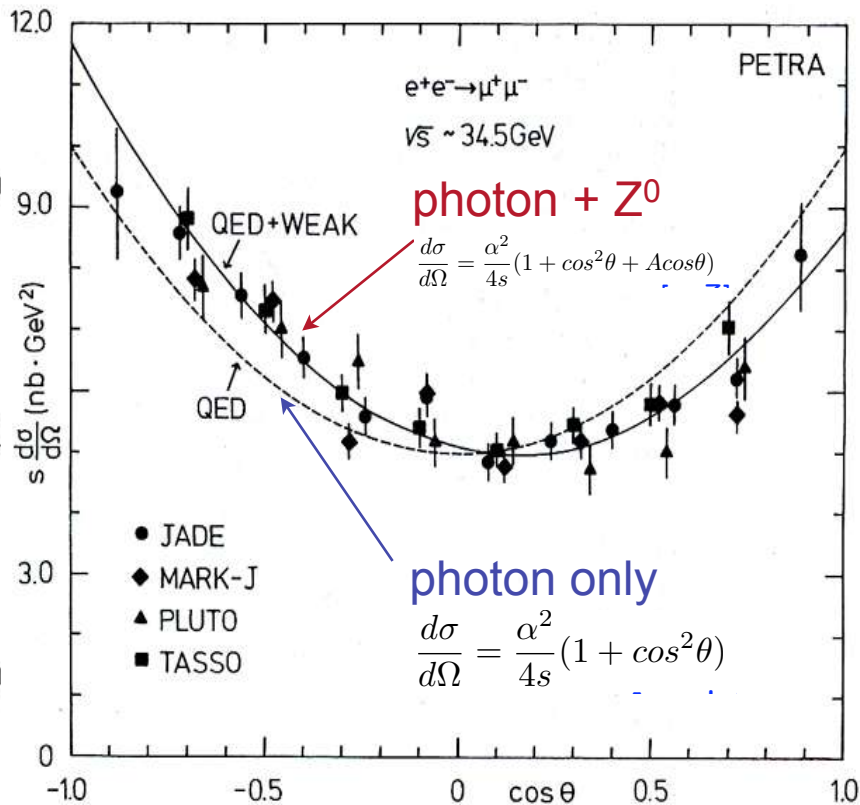


Example: event display

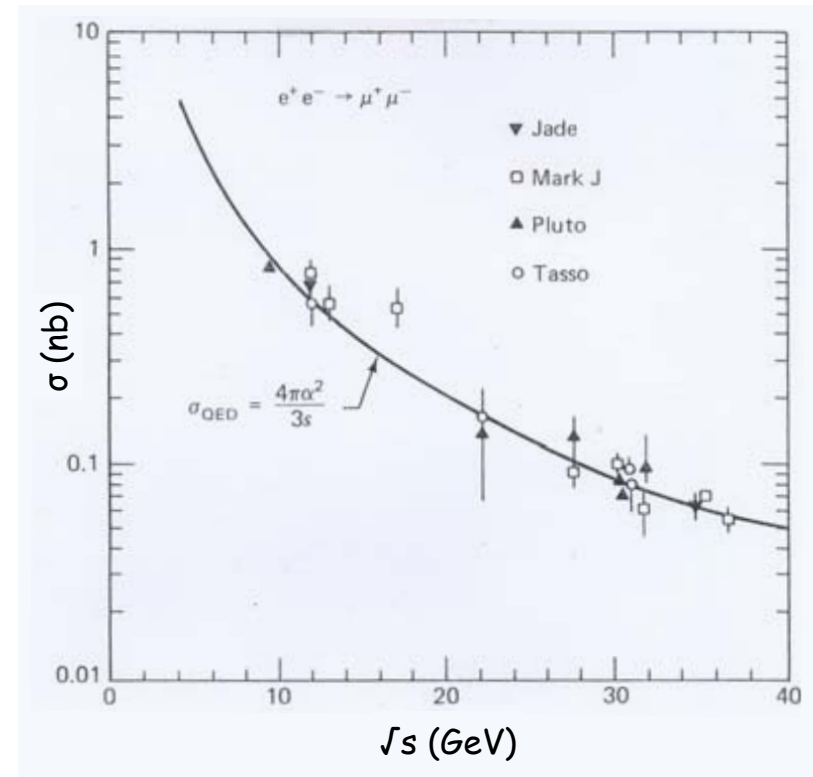


Muon pair production

Differential cross section

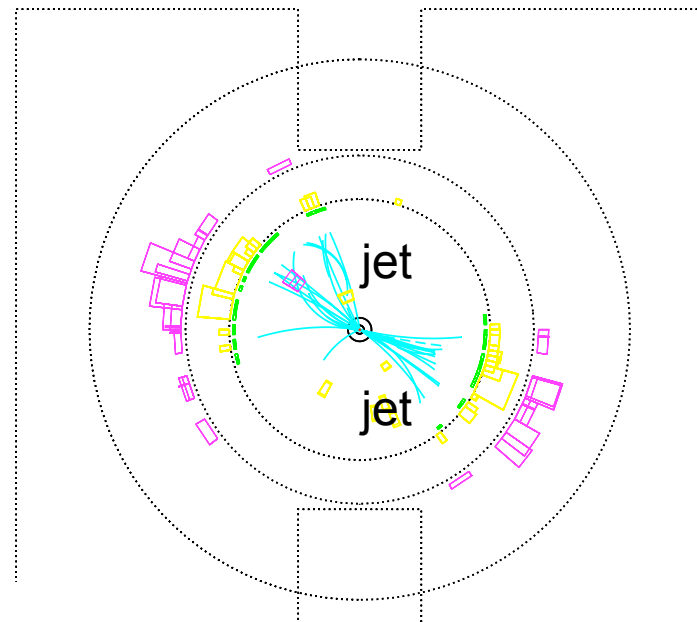
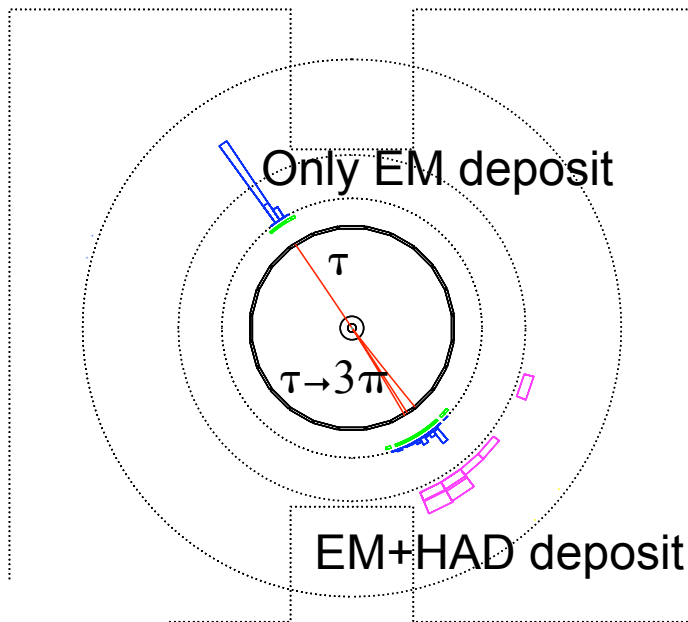
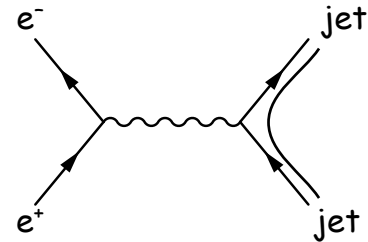
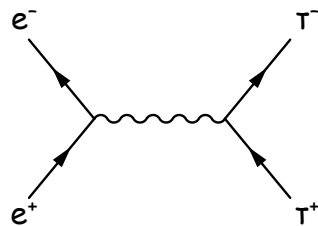


Total cross section vs. energy



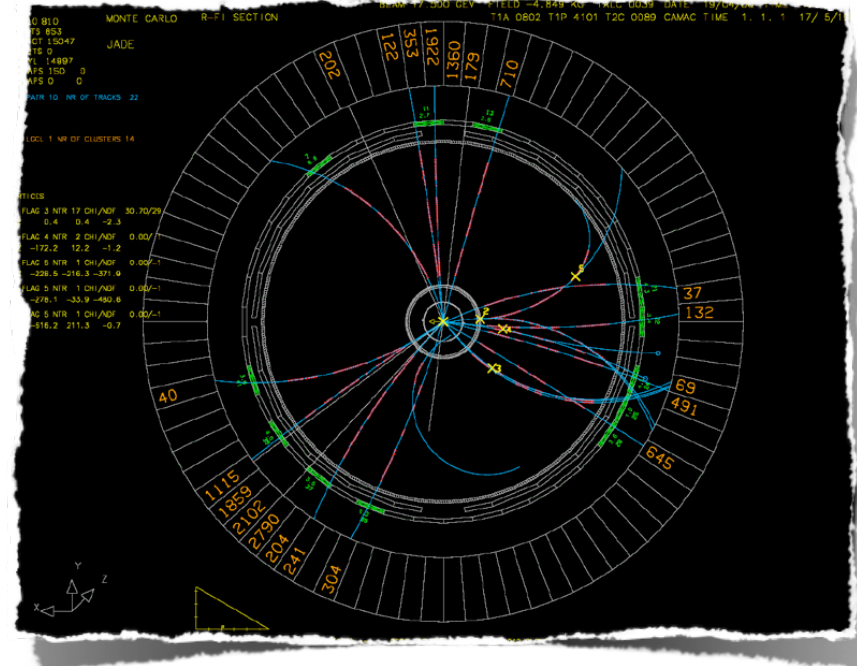
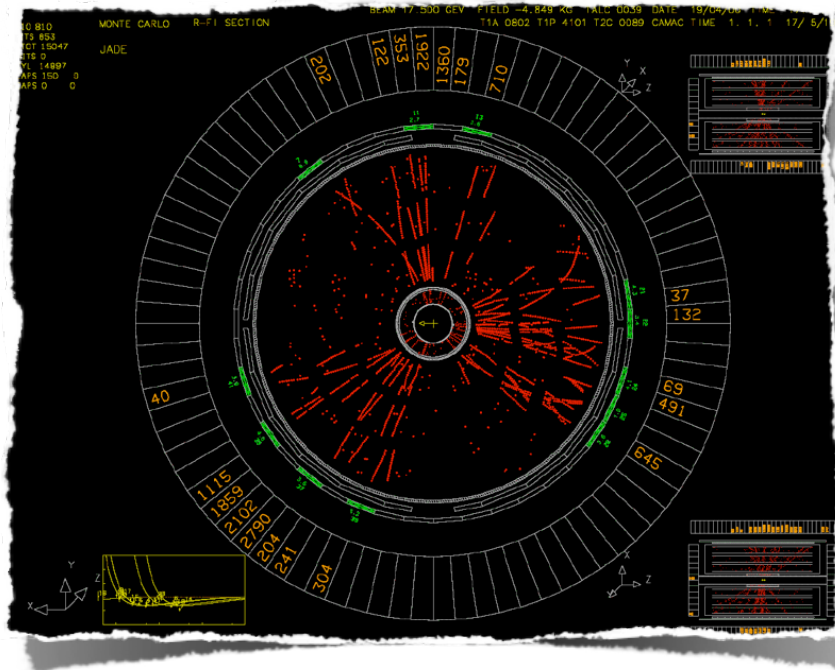
- Angular distribution is sensitive to electroweak corrections due to Z^0 exchange
 - ◆ Additional term proportional to $\cos(\theta)$
- Total cross section (integrating over solid angle) goes as $1/s$

Other final states



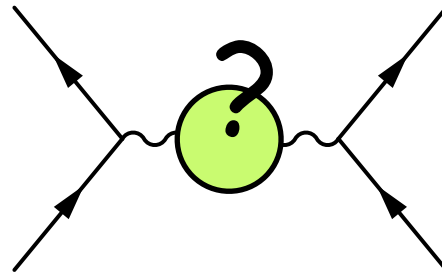
3 jets final states

Hard gluon radiation from final quark-antiquark pair



Limits of QED

Question: what do we expect if QED is not the only physics involved in the scattering processes discussed so far?



We define an energy scale Λ , or distance $r \sim 1/\Lambda$, at which the QED theoretical model does not describe the data

Change of potential

$$\frac{1}{r} \rightarrow \frac{1}{r} (1 - e^{-\Lambda r})$$

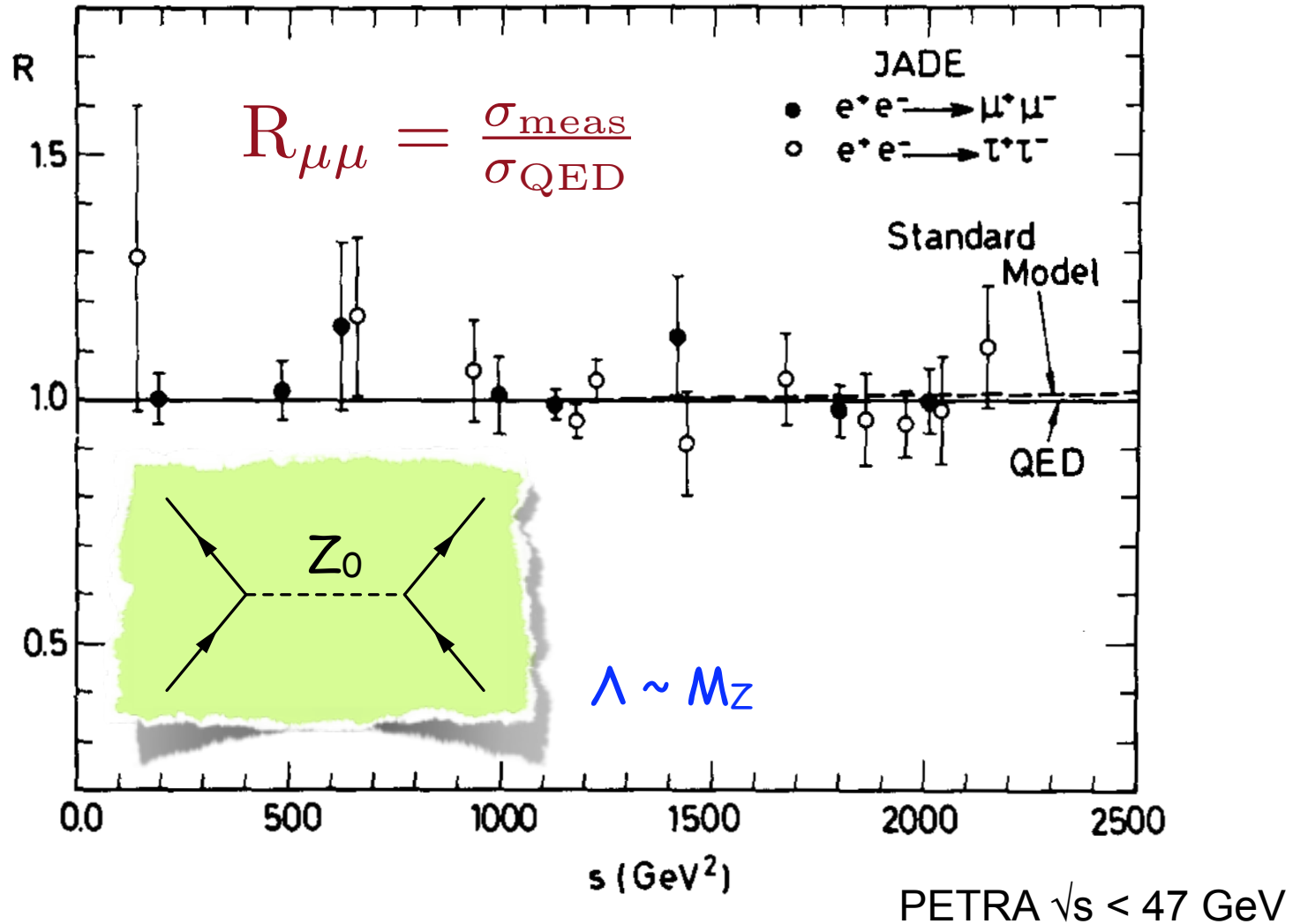
Change of propagator

$$-\frac{1}{q^2} \rightarrow -\frac{1}{q^2} \left(1 + \frac{q^2}{\Lambda^2} \right)$$

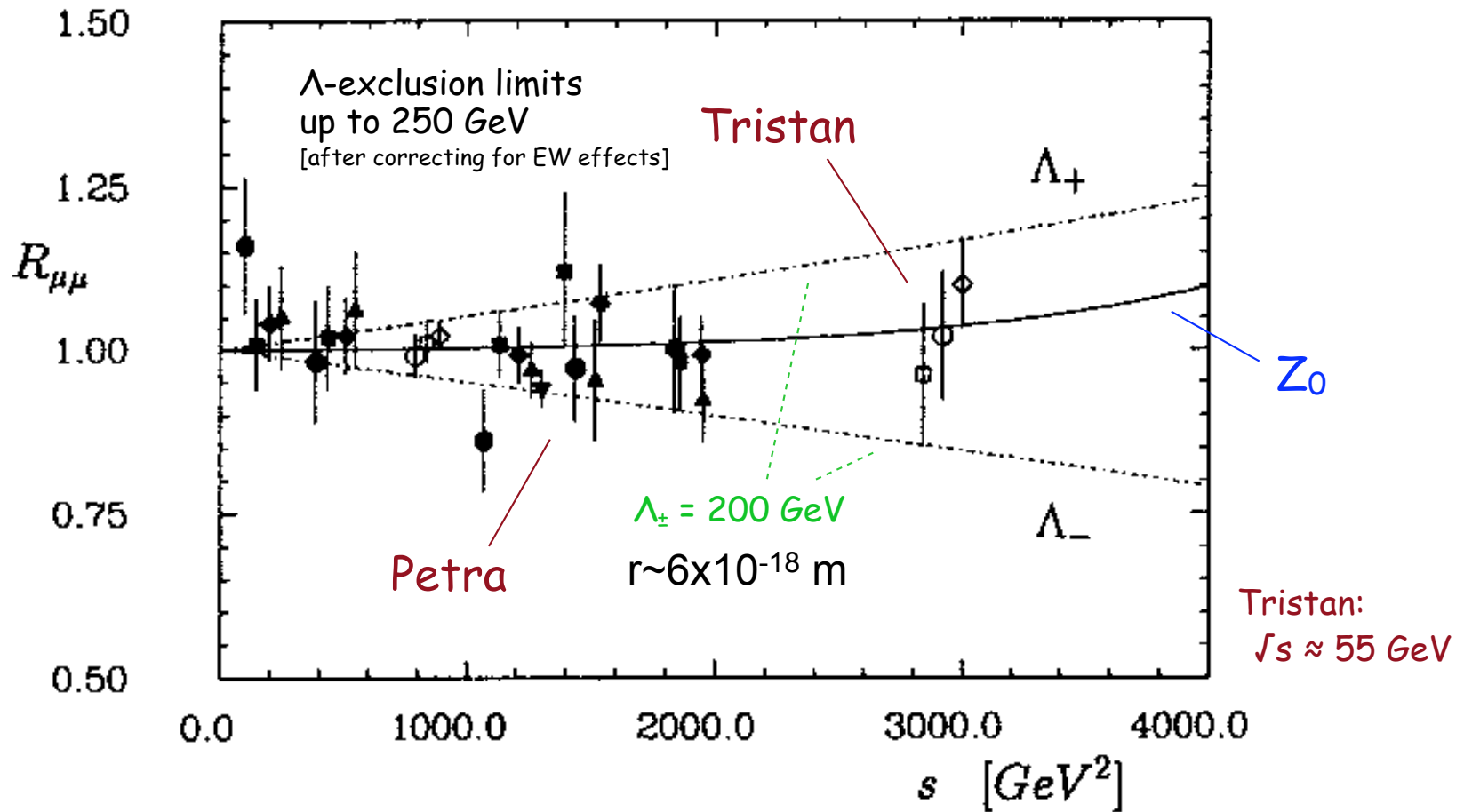
Change of cross section

$$\sigma_{\mu\mu} = \frac{3\pi\alpha^2}{3s} \left(1 \pm \frac{s}{\Lambda^2 - s} \right)^2$$

Beyond QED



Beyond QED



Part two:

*Measurement of the electron
anomalous magnetic moment*

Electron magnetic moment

- Rotating electrically charged body creates a magnetic dipole
 - ◆ External magnetic field exerts a torque on the electron magnetic moment
- Electrons have intrinsic **magnetic moment**, related to their spin

$$\mu = -g \frac{e}{2m} s = -\frac{g}{2} \frac{e}{2m}$$

- In case of electrons the magnetic moment is **anti-parallel** to the spin
- The **g-factor is equal to 2**, as calculated from Dirac's equation

$$a = \frac{\alpha}{2\pi} = \frac{g - 2}{2} = 0$$

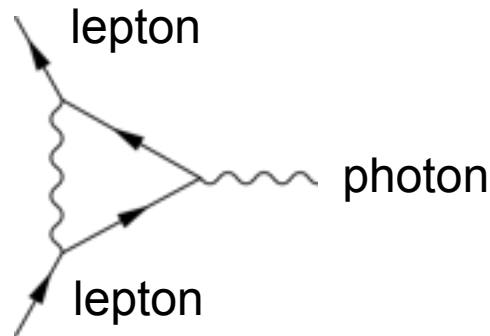
- Corrections to the g-factor are given by **higher order QED contributions** as well as **hadronic** and **weak interactions**. There could be additional contributions from **physics beyond the Standard Model (SM)**

$$\frac{g}{2} = 1 + a_{QED}(\alpha) + a_{hadronic} + a_{weak} + a_{new}$$

- When adding the corrections we usually talk of **anomalous magnetic moment of the electron**

QED: higher order corrections

- The one-loop corrections to the magnetic moment are due to vacuum fluctuation and polarization effects. For example:

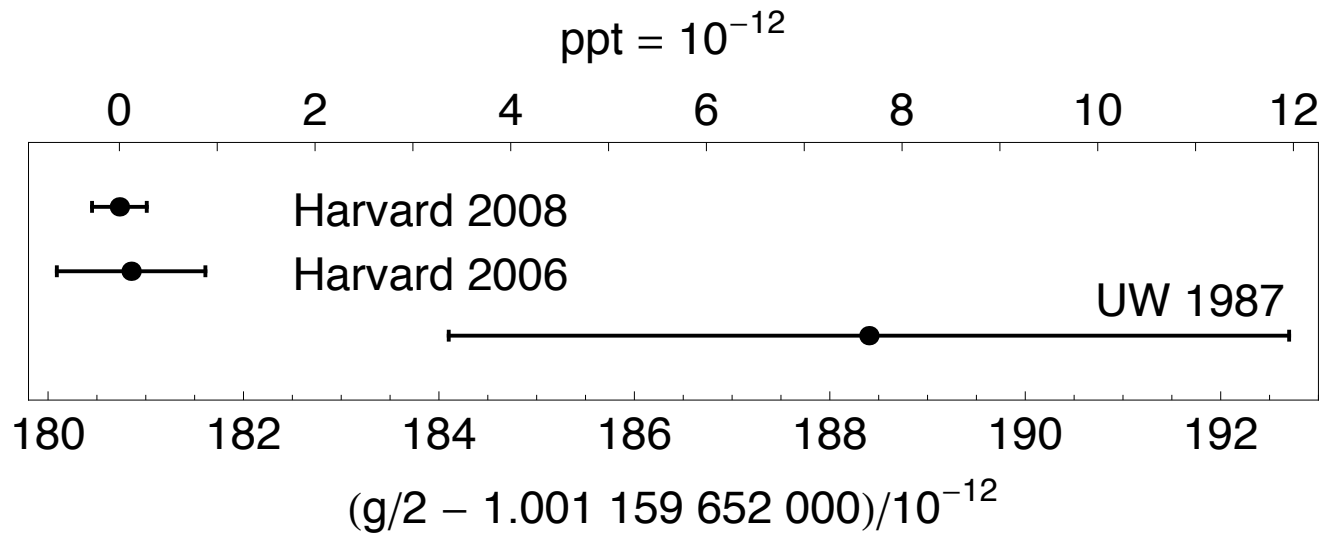


- The textbook calculation of the one-loop corrections gives corrections $\sim 10^{-3}$ (see References):

$$a = \frac{\alpha}{2\pi} = \frac{g - 2}{2} \simeq .0011614$$

- Hadronic and weak interactions are calculated (within the SM) to be very small and negligible, respectively
- As we will see, the precision achieved by experimental results need QED predictions with α^4 **precision**

Current status of $g/2$ measurements



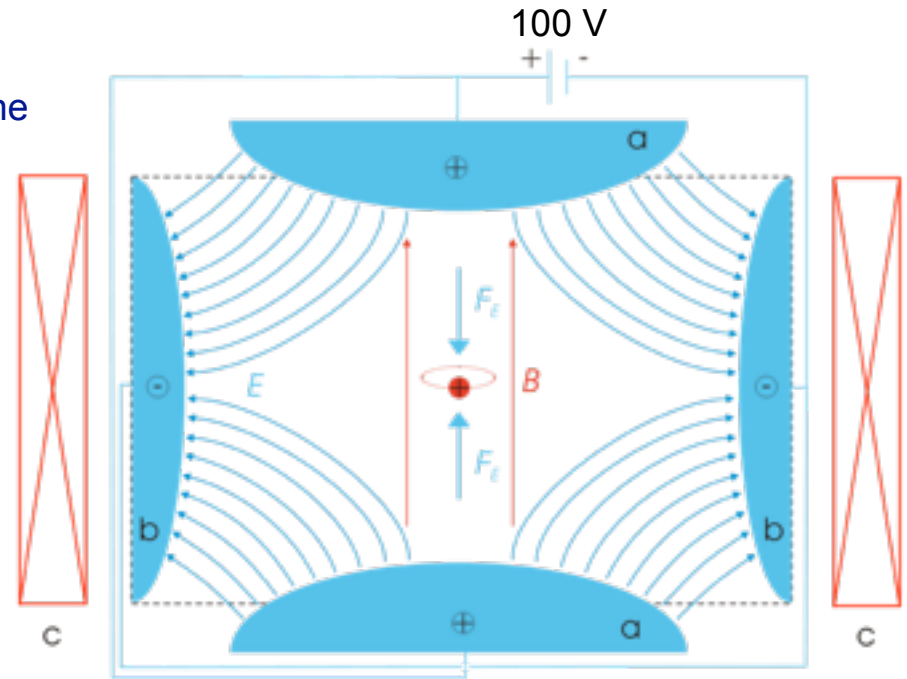
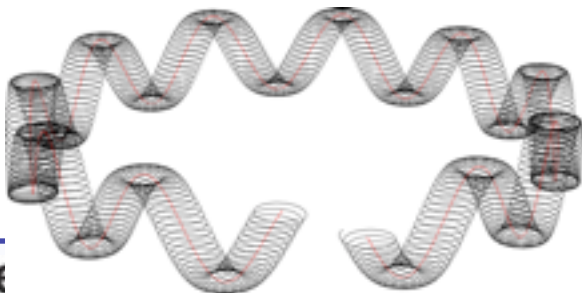
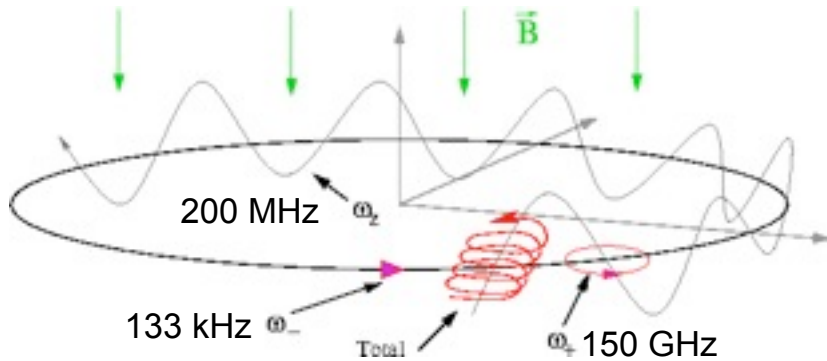
- The precision nowadays is below 10^{-12} !
- Latest measurements 15 times more precise than previous result, which stood for about 20 years
- Measured value is shifted by 1.7 standard deviations
- How did we get to this astonishing precision?

Experiment: main ingredients

- Single-electron quantum cyclotron:
 - ◆ *A Penning trap suspends and confines the electron in an atom-like state*
- Fully resolved cyclotron and spin energy levels:
 - ◆ *Accurate measurements of the resonant frequencies of driven transitions between the energy levels of this homemade atom – an electron bound to our trap – reveals the electron magnetic moment in units of Bohr magnetons, $g/2$*
- Detection sensitivity sufficient to detect one quantum transitions
 - ◆ *Frequency detection sensitivity in the radio and microwave region*

Penning trap

- Penning trap confines electrons by using:
 - ◆ A strong **vertical magnetic field** to confine particles radially
 - ◆ A **quadrupole electric field** to confine particles axially
- The magnetic field is produced by a **solenoid**
- The electric field is produced by three electrodes: one **ring** and two **endcaps**
- The trajectory in the radial plane is characterized by two frequencies



Magnetron frequency: ω_-
 Modified cyclotron frequency: ω_+
 The cyclotron frequency is $(\omega_+ + \omega_-)$
 A small-frequency oscillation is also in the vertical plane (axial frequency ω_z)

Energy levels measurement

- A non-relativistic electron in a magnetic field has energy levels:

$$E(n, m_s) = \frac{g}{2} h \nu_c m_s + \left(n + \frac{1}{2} \right) h \nu_c \quad \nu_c = \frac{eB}{2\pi m} \quad \nu_s = \frac{g}{2} \nu_c = \frac{g}{2} \frac{eB}{2\pi m}$$

- Depend on the **cyclotron frequency** (ν_c) and on the **spin frequency** (ν_s)

$$\frac{g}{2} = \frac{\nu_s}{\nu_c} = 1 + \frac{\nu_s - \nu_c}{\nu_c} = 1 + \frac{\nu_a}{\nu_c}$$

- Since ν_s and ν_c differ only by a part per 10^3 measuring ν_a and ν_c to a precision of 1 part on 10^{10} gives $g/2$ to **1 part to 10^{13}** .
- Two advantages of this technique:
 - ◆ One can measure the ratio of two frequency to very high precision
 - ◆ Since the B field appears in both numerator and denominator, the dependence on the magnetic field cancels in the ratio

Energy levels measurement

- Including the relativistic corrections, the energy levels are given by:

$$E(n, m_s) = \frac{g}{2} h \nu_c m_s + (n + \frac{1}{2}) h \bar{\nu}_c - \frac{1}{2} h \delta (n + \frac{1}{2} + m_s)^2$$

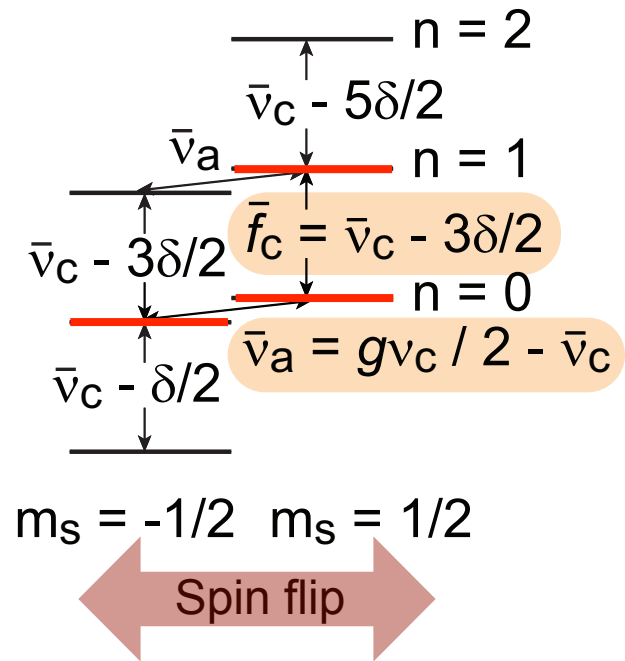
Relativistic correction term
 ↓
Frequency shift due to Penning trap

- The experiment measures the following transition frequencies:

$$\bar{f}_c \equiv \bar{\nu}_c - \frac{3}{2} \delta \quad 1, 1/2 \rightarrow 0, 1/2$$

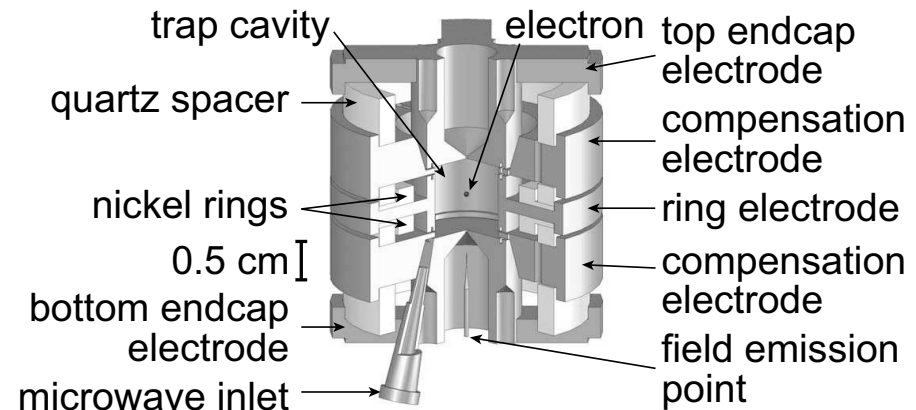
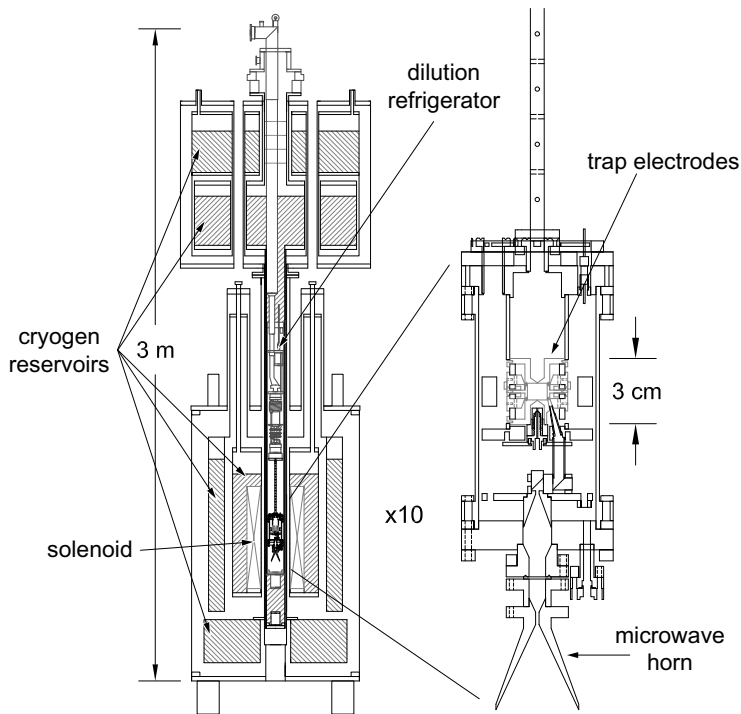
$$\bar{\nu}_a \equiv \frac{g}{2} \nu_c - \bar{\nu}_c \quad 0, 1/2 \rightarrow 0, -1/2$$

Cyclotron frequency ~ 150 GHz



Experimental setup

- A Penning trap is used to artificially bound the electron in an orbital state
- High voltage (100V) applied between cylindric and endcap contacts

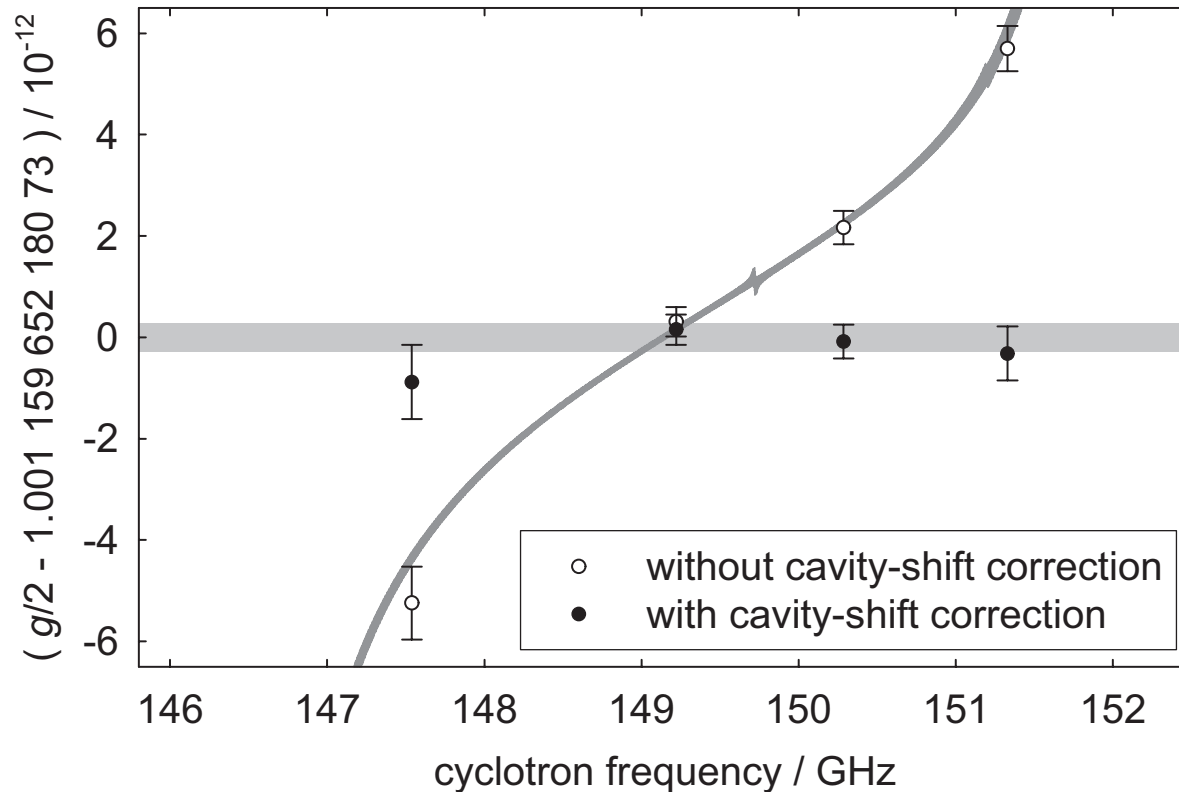


- A high magnetic field (5 T) is necessary to increase the spacing between cyclotron energy levels ($\nu_c \propto B$)
- Very low temperature (100 mK) increases the probability to populate the orbital ground state

$$P \propto \exp[-h\bar{\nu}_c/(kT)]$$

Results

$$g/2 = 1.001\,159\,652\,180\,73\,(28) \quad [0.28 \text{ ppt}]$$



Shifts are induced by interaction of electron
with nearby cavity radiation modes

Solution: do measurements at various frequencies

Theoretical predictions

- The QED calculations provide the prediction for $g/2$ up to the fourth power of alpha:

$$\frac{g}{2} = 1 + C_2 \left(\frac{\alpha}{\pi}\right) + C_4 \left(\frac{\alpha}{\pi}\right)^2 + C_6 \left(\frac{\alpha}{\pi}\right)^3 + C_8 \left(\frac{\alpha}{\pi}\right)^4 + C_{10} \left(\frac{\alpha}{\pi}\right)^5 + \dots + a_{\text{hadronic}} + a_{\text{weak}},$$
$$C_2 = 0.500\,000\,000\,000\,00 \text{ (exact)}$$
$$C_4 = -0.328\,478\,444\,002\,90 \text{ (60)}$$
$$C_6 = 1.181\,234\,016\,827 \text{ (19)}$$
$$C_8 = -1.914\,4 \text{ (35)}$$
$$C_{10} = 0.0 \text{ (4.6).}$$
$$a_e^{\text{hadronic}} = 1.682(20) \times 10^{-12}$$

- From this formula and theoretical predictions we can:

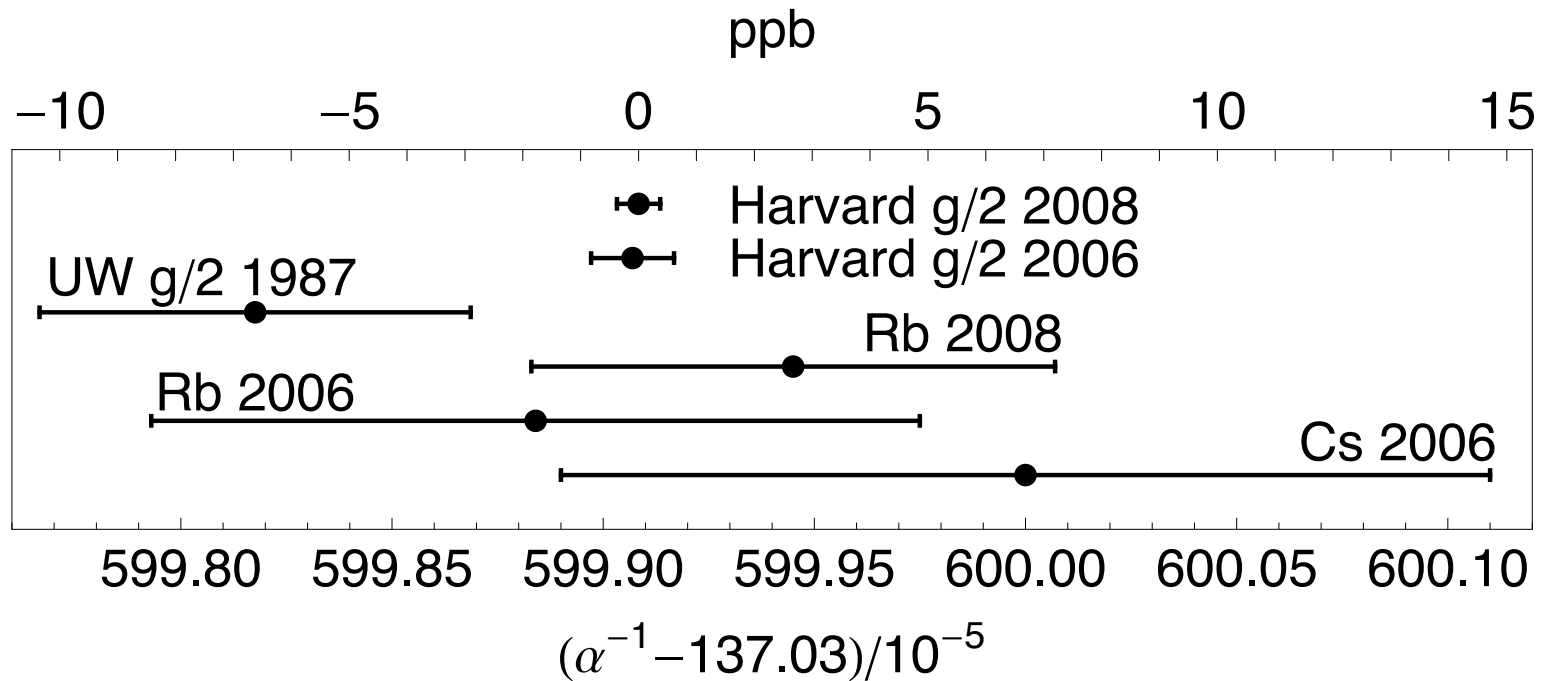
- Measure the coupling constant α

$$\alpha^{-1} = 137.035\,999\,084 \text{ (33) (39) [0.24 ppb] [0.28 ppb]}$$
$$= 137.035\,999\,084 \text{ (51) [0.37 ppb].}$$

- Comparing the measured $g/2$ with expectation using α from other measurements

$$g/2 = 1.001\,159\,652\,180\,73 \text{ (28) [0.28 ppt], } \textit{Measured}$$
$$g(\alpha)/2 = 1.001\,159\,652\,177\,60 \text{ (520) [5.2 ppt]. } \textit{Predicted}$$

Status of high-precision α measurements



Source: <http://hussle.harvard.edu/~gabrielse/gabrielse/papers/2009/DeterminingTheFineStructureConstant.pdf>

References

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 - M.Peskin, D.Schroeder, *An Introduction to Quantum Field Theory*, Addison-Wesley, 1995
- Measurement of the electron magnetic moment:
 - G.Gabrielse, *Lepton Dipole Moments: The Search for Physics Beyond the Standard Model*, edited by B.L. Roberts and W.J. Marciano (World Scientific, Singapore, 2009), Advanced Series on Directions in High Energy Physics – Vol. 20.
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 - <http://hussle.harvard.edu/~gabrielse/gabrielse/papers/2009/DeterminingTheFineStructureConstant.pdf>
- Experimental QED tests at e+e- colliders:
 - A.Ali, P.Soeding, *High Energy Electron Positron Physics*, World Scientific
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