

Exercise 10.1 Teleportation Redux

- (a) Show that for the entangled state $|\Phi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ and any unitary operator U ,

$$(U_A \otimes \bar{U}_B) |\Phi\rangle_{AB} = |\Phi\rangle_{AB},$$

where $\bar{}$ denotes complex conjugation in the $|0\rangle, |1\rangle$ basis.

- (b) Show that for any state $|\psi\rangle$

$${}_A\langle\psi|\Phi\rangle_{AB} = \frac{1}{\sqrt{2}}|\psi^*\rangle_B.$$

- (c) Use the results of (a) and (b) to give a derivation of the teleportation protocol without resorting to components.
- (d) What happens if Alice and Bob use the state $(\mathbb{1}_A \otimes U_B)|\Phi\rangle_{AB}$ for teleportation? Or if Alice measures in the basis $\bar{U}_{A'}|\Phi_j\rangle_{A'A}$?
- (e) Instead of a single system state $|\psi\rangle_{A'}$, Alice has a bipartite state $|\psi\rangle_{A_1A_2}$. What happens if she performs the teleportation protocol on system A_2 ?

Exercise 10.2 Remote Copy

Alice and Bob would like to create the state $|\Psi\rangle_{AB} = a|00\rangle_{AB} + b|11\rangle_{AB}$ from Alice's state $|\psi\rangle_A = a|0\rangle_A + b|1\rangle_A$, a “copy” in the quantum-mechanical sense. Additionally, they share the canonical entangled state $|\Phi\rangle$. Can they create the desired state by performing only local operations (measurements and unitary operators), provided Alice can only send *one* bit of classical information to Bob?

Exercise 10.3 Quantum mutual information

Consider a composed system $A \otimes B \otimes C$ with a shared state ρ_{ABC} .

In a first step we ignore system C and consider only $A \otimes B$ (and the reduced state $\rho_{AB} = \text{Tr}_C(\rho_{ABC})$). One way of quantifying the correlations between A and B is to use the *mutual information* between them, defined as

$$I(A : B) = H(A) + H(B) - H(AB) \quad (1)$$

$$= H(A) - H(A|B). \quad (2)$$

If we have access to C , we can define a conditional version of the mutual information between A and B as

$$I(A : B|C) = H(A|C) + H(B|C) - H(AB|C) \quad (3)$$

$$= H(A|C) - H(A|BC). \quad (4)$$

- (a) Assume a system formed by two qubits A and B that share a state ρ_{AB} . Consider bases $\{|0\rangle_A, |1\rangle_A\}$ and $\{|0\rangle_B, |1\rangle_B\}$ for the subsystems of each qubit.

1. Check that the mutual information of the fully entangled state, $|\Psi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$, is maximal.

2. See that for classically correlated states, $\rho_{AB} = p|0\rangle\langle 0|_A \otimes \sigma_B^0 + (1-p)|1\rangle\langle 1|_A \otimes \sigma_B^1$ (where $0 \leq p \leq 1$), the mutual information cannot be greater than one.

(b) Consider the so-called *cat state* shared by four qubits, $A \otimes B \otimes C \otimes D$, that is defined as

$$|\ominus\rangle = \frac{1}{\sqrt{2}} (|0000\rangle + |1111\rangle). \quad (5)$$

Check how the mutual information between qubits A and B changes with the knowledge of the remaining qubits, namely:

1. $I(A : B) = 1$.
2. $I(A : B|C) = 0$.
3. $I(A : B|CD) = 1$.